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Measurement of the Formation Rate of Muonic Hydrogen Molecules

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I. INTRODUCTION

Nuclear muon capture on the proton,

\[ \mu^- + p \rightarrow n + v_\mu, \]  \hspace{1cm} (1)

is a basic charged-current weak reaction [1–3]. Several experiments have measured the rate of ordinary muon capture [Eq. (1)] or the rarer process of radiative muon capture, \( \mu + p \rightarrow n + v + \gamma \), to determine the weak pseudoscalar coupling of the proton, \( g_P \), which can be extracted most straightforwardly from muon capture on the nucleon. A precision determination of \( g_P \) has been a longstanding experimental challenge [2,3] due to the small rate of capture on the proton and complications arising from the formation of muonic molecules. The most recent MuCap result, \( g_P = 8.06 \pm 0.55 \) [4], achieved an unprecedented precision of 7\%, thereby providing a sensitive test of QCD symmetries and confirming a fundamental prediction of chiral perturbation theory, \( g_P = 8.26 \pm 0.23 \) [5–7].

Experimentally, process (1) is observed after low-energy muons are stopped in hydrogen, where they form \( p\mu \) atoms and \( p\mu\mu \) molecules. The overlap in the wave functions of
the proton and the bound muon leads to small but observable capture rates at the $10^{-3}$ level relative to muon decay, $\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$, which is the dominant mode of muon disappearance in that environment. The nuclear capture rates depend on the spin compositions of the muonic atoms and molecules (a direct consequence of the V-A structure of the electroweak interaction), and thus the rates vary significantly among the different muonic states. The calculated rates for the two hyperfine states of the $p\mu$ atom possessing spin $F = 0$ and $1$ are $\lambda_S = 712.7$ s$^{-1}$ and $\lambda_T = 12.0$ s$^{-1}$, respectively (cf. Ref. [1], updated in Ref. [4]). The formation of $p\mu\mu$ molecules further complicates the situation, as the calculated capture rates for the ortho and para states, $\lambda_{OM} = 542.4$ s$^{-1}$ and $\lambda_{PM} = 213.9$ s$^{-1}$, differ too from the atomic rates [Eq. (4)]. Correct interpretation of the observed muon disappearance rate thus relies on a thorough understanding of the “muon chemistry” reactions governing the time evolution of the $p\mu$ and $p\mu\mu$ states. This interrelationship between muon capture and muon chemistry in hydrogen has been the primary source of ambiguity in the 50-year history of experiments in the field. Historically, interest in muon atomic and molecular reactions arose due to their abovementioned relevance for the determination of nuclear muon capture rates in hydrogen isotopes [2] and their importance in muon-catalyzed fusion [8], where $\lambda_{pp\mu}$ was calculated within a systematic program to solve the Coulomb three-body problem [9].

The MuCap experiment employed a novel technique involving the use of low-pressure hydrogen gas to suppress molecular formation. Nevertheless, it was still necessary to apply corrections that were based on measurements of the molecular formation rates that determine the $p\mu\mu$ ortho and para molecule populations. In the initial MuCap physics result [10], we conservatively estimated that the uncertainty in the molecular formation rate $\lambda_{pp\mu}$ contributed a systematic uncertainty of $4.3$ s$^{-1}$ to our determination of $\Lambda_S$, the muon capture rate in the $p\mu$ hyperfine singlet state. During the later high statistics data taking for MuCap, we performed a dedicated measurement of $\lambda_{pp\mu}$ to improve the precision on this parameter and render its contribution to the uncertainty on $\Lambda_S$ nearly negligible. The final MuCap result, $\Lambda_S = (714.9 \pm 5.4_{\text{stat}} \pm 5.1_{\text{syst}})$ s$^{-1}$ [4], possessed greatly improved statistical and systematic uncertainties. A preliminary value for $\lambda_{pp\mu}$ obtained from our measurement was an important ingredient in this result. In this paper we document the $\lambda_{pp\mu}$ experiment and present its final results.

The contents of this article are as follows. In Sec. II we introduce muon-induced processes in hydrogen and their impact on muon-capture measurements. In Sec. III we describe the MuCap experiment and our technique for measuring $\lambda_{pp\mu}$; the corresponding data analysis and result for $\lambda_{pp\mu}$ are described in Sec. IV. In Sec. V we use our new result for $\lambda_{pp\mu}$ to update previous MuCap measurements. A concluding summary is given in Sec. VI.

II. MUON CAPTURE AND MUON CHEMISTRY

A. Muon reactions in hydrogen

Muons stopped in hydrogen can form a variety of atomic and molecular states that are subject to different physical processes and whose populations are governed by the rates shown in Fig. 1. Table I lists all of the rates used in this paper and their values. Several of the atomic processes proceed via binary collisions of muonic atoms with other target molecules. It is conventional to normalize those density-dependent rates to the values observed at the LH$_2$ density, $\phi_0 = 4.25 \times 10^{22}$ atoms/cm$^3$, and express all target densities $\phi$ relative to $\phi_0$.

![Figure 1](image)

**FIG. 1.** (Color online) Kinetics of negative muons in hydrogen containing a single chemical impurity $Z$. The large circles represent the different muonic states that can form. The black arrows denote transitions between muonic states, while the colored arrows with a small circle at the beginning of the line indicate muon disappearance due to the weak-interaction processes of decay or nuclear capture. The arrow labeled $\lambda_S$ is dashed to indicate that its rate is about 240 times smaller than that of $\lambda_{cd}$.

<table>
<thead>
<tr>
<th>Muon process</th>
<th>Symbol</th>
<th>Value (s$^{-1}$)</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Weak-interaction rates</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Muon decay</td>
<td>$\lambda_-$</td>
<td>455170.05 $\pm$ 0.46</td>
<td>[11–13]</td>
</tr>
<tr>
<td>$p\mu$ singlet capture</td>
<td>$\Lambda_S$</td>
<td>714.9 $\pm$ 7.4</td>
<td>[4]</td>
</tr>
<tr>
<td>$p\mu$ triplet capture</td>
<td>$\Lambda_T$</td>
<td>12.0 $\pm$ 0.1</td>
<td>[1]</td>
</tr>
<tr>
<td>$p\mu\mu$ ortho capture</td>
<td>$\Lambda_{OM}$</td>
<td>544.0 $\pm$ 11.3</td>
<td>Eq. (4)</td>
</tr>
<tr>
<td>$p\mu\mu$ para capture</td>
<td>$\Lambda_{PM}$</td>
<td>214.6 $\pm$ 4.2</td>
<td>Eq. (4)</td>
</tr>
<tr>
<td>N capture</td>
<td>$\Lambda_N$</td>
<td>6.93 $\pm$ 0.08 $\times$ 10$^4$</td>
<td>[14]</td>
</tr>
<tr>
<td>O capture</td>
<td>$\Lambda_O$</td>
<td>10.26 $\pm$ 0.06 $\times$ 10$^4$</td>
<td>[14]</td>
</tr>
<tr>
<td>Ar capture</td>
<td>$\Lambda_{Ar}$</td>
<td>141 $\pm$ 11 $\times$ 10$^4$</td>
<td>[14]</td>
</tr>
<tr>
<td>&amp;</td>
<td>130.2 $\pm$ 3.2 $\times$ 10$^4$</td>
<td>this work</td>
<td></td>
</tr>
<tr>
<td>Atomic and molecular rates</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>o-p transition</td>
<td>$\lambda_{op}$</td>
<td>6.6 $\pm$ 3.4 $\times$ 10$^4$</td>
<td>[2]</td>
</tr>
<tr>
<td>$p\mu\mu$ formation$^a$</td>
<td>$\lambda_{pp\mu}$</td>
<td>2.3 $\pm$ 0.2 $\times$ 10$^6$</td>
<td>[2]</td>
</tr>
<tr>
<td>Transfer to N$^a$</td>
<td>$\lambda_{pN}$</td>
<td>0.34 $\pm$ 0.07 $\times$ 10$^{11}$</td>
<td>[15]</td>
</tr>
<tr>
<td>Transfer to O$^a$</td>
<td>$\lambda_{pO}$</td>
<td>0.85 $\pm$ 0.02 $\times$ 10$^{11}$</td>
<td>[16]</td>
</tr>
<tr>
<td>Transfer to Ar$^a$</td>
<td>$\lambda_{pAr}$</td>
<td>1.63 $\pm$ 0.09 $\times$ 10$^{11}$</td>
<td>[17]</td>
</tr>
<tr>
<td>&amp;</td>
<td>1.94 $\pm$ 0.11 $\times$ 10$^{11}$</td>
<td>this work</td>
<td></td>
</tr>
<tr>
<td>Ar Huff factor$^b$</td>
<td>$h$</td>
<td>0.985 $\pm$ 0.003</td>
<td>[18–20]</td>
</tr>
</tbody>
</table>

$^a$Normalized to LH$_2$ density $\phi_0$.
$^b$Dimensionless quantity.
Negatively charged, low-energy muons entering hydrogen are slowed down and undergo atomic capture, forming highly excited \( p\mu \) atoms. After an atomic cascade to the ground state, the two hyperfine states of the \( p\mu \) atom, singlet \( (F = 0) \) and triplet \( (F = 1) \), are populated according to their statistical weights, \( \frac{1}{2} \) and \( \frac{3}{2} \), respectively. These complex initial stages happen on a time scale of nanoseconds at target densities exceeding \( \phi \gtrsim 0.01 \), as in our case. Charge-exchange collisions \([21,22]\) convert the higher-lying triplet state to the lower-lying singlet state at a rate calculated to be \( \approx \phi \times 2 \times 10^{10} \text{s}^{-1} \) \([23]\). Thus after less than 100 ns the triplet state is effectively depopulated and the main features of the kinetics can be described by the scheme depicted in Fig. 1. This condition is true for the present analysis and in Refs. \([4,10]\).

For our purposes, muon kinetics in pure hydrogen effectively starts with the \( p\mu \) atom in its hyperfine singlet state. In subsequent collisions of the \( p\mu \) atom with hydrogen molecules, two types of \( p\mu \) molecules can be formed that differ in their angular momentum \( L \) and total spin \( I \). Due to the Fermi statistics of the two-proton system, the ortho state \((pp\mu)_\text{om}\) has \( L = 1 \) and \( I = 1 \), while the para state \((pp\mu)_\text{pm}\) has \( L = 0 \) and \( I = 0 \). According to theory, \( pp\mu \) formation proceeds to the ortho state predominantly at the normalized rate \( \lambda_{\text{of}} = 1.8 \times 10^{9} \text{s}^{-1} \), while the para formation rate \( \lambda_{\text{pf}} = 7.5 \times 10^{3} \text{s}^{-1} \) is much smaller \([9]\). The total normalized molecular formation rate is the sum of these two rates:

\[
\lambda_{pp\mu} = \lambda_{\text{of}} + \lambda_{\text{pf}}.
\]

Molecular formation scales with the target density \( \phi \), so experiments observe the effective molecular formation rate

\[
\Lambda_{pp\mu} = \phi \lambda_{pp\mu}.
\]

The transition from the \( pp\mu \) ortho state to the para state at rate \( \lambda_{\text{op}} \) involves a proton spin flip and is only allowed due to relativistic effects in the molecular wave function. The \( (pp\mu_{\text{om}})_{\text{p}} \) is positively charged and quickly forms various molecular complexes in collisions with \( \text{H}_2 \) molecules. The ortho-para transition proceeds at the calculated rate \( \lambda_{\text{op}} = (7.1 \pm 1.2) \times 10^{5} \text{s}^{-1} \) via the emission of an electron from these clusters \([24]\). Two previous experiments measured \( \lambda_{\text{op}} \) and obtained the inconsistent results \((4.1 \pm 1.4) \times 10^{4} \text{s}^{-1}\) \([25]\) and \((11.1 \pm 1.9) \times 10^{4} \text{s}^{-1}\) \([26]\). Review \([2]\) therefore inflated the uncertainties and quoted an average experimental value of \( \lambda_{\text{op}} = (6.6 \pm 3.4) \times 10^{4} \text{s}^{-1} \), which we use in this work.

As mentioned above, the weak nuclear capture rates strongly depend on spin factors within the total \( pp\mu \) molecular spin function and can be expressed as

\[
\Lambda_{\text{OM}} = 2\gamma_{\text{om}} \left( \frac{3}{8} \Lambda_{S} + \frac{1}{8} \Lambda_{T} \right),
\]

\[
\Lambda_{\text{PM}} = 2\gamma_{\text{pm}} \left( \frac{5}{8} \Lambda_{S} + \frac{3}{8} \Lambda_{T} \right).
\]

The molecular overlap factors are \( 2\gamma_{\text{om}} = 1.009 \pm 0.001 \) and \( 2\gamma_{\text{pm}} = 1.143 \pm 0.001 \) \([24]\). Based on these equations the capture rates of the molecular states can be calculated using the MuCap result for \( \Lambda_{S} \) and the theoretical value for the smaller rate \( \Lambda_{T} \) as input (see Table I).

In the presence of \( Z > 1 \) chemical impurities, the muon can form a bound \( Z\mu \) state instead of a \( p\mu \) atom. The factor \( f \) in Fig. 1 characterizes the initial population of \( Z\mu \) atoms, which arises from two pathways. First, at the time of the muon stop, \( Z > 1 \) elements are energetically favored over hydrogen by Coulomb capture. Second, during the \( p\mu \) deexcitation cascade, prompt transfer to higher \( Z \) elements can occur. The size of \( f \) scales linearly with the relative atomic concentration \( c_{Z} \) of the impurity.

Muons will also transfer from the singlet \( p\mu \) state to the energetically favorable \( Z\mu \) state in collisional processes. Transfer from the molecular states to the \( Z\mu \) state is not possible because the charged \((pp\mu)^{+}\) molecule is repelled by the \( Z \) nucleus. The effective transfer rate to the impurity \( \Lambda_{pZ} \) is expressed as

\[
\Lambda_{pZ} = c_{Z} \phi \lambda_{pZ},
\]

where \( \lambda_{pZ} \) is the normalized transfer rate. Excited \( Z\mu \) states are created by such transfers, and observable muonic x rays are emitted during the subsequent deexcitation cascade. The rate \( \Lambda_{Z} \) of subsequent muon capture on the nucleus increases roughly proportional to \( Z^{4} \) (the more realistic Primakoff formula is discussed in Ref. \([14]\)). Table I shows that the capture rates for typical impurity elements (nitrogen, oxygen, and argon) are all much larger than the \( p\mu \) singlet capture rate \( \Lambda_{S} \).

The natural abundance of deuterium in hydrogen generally causes an additional loss channel due to the formation of \( d\mu \) atoms \([27,28]\) and \( pd\mu \) molecules \([3]\). For the presented measurement, a cryogenic distillation column was used to isotopically purify the hydrogen, achieving a final deuterium concentration of less than 10 ppb \([4]\). At this level, the deuterium loss channel is completely negligible.

The muon can decay from any of the states in Fig. 1 at a rate close to the free muon decay rate \( \lambda_{\mu} \) \([11,13]\). The actual decay rates are slightly reduced with respect to \( \lambda_{\mu} \) by the Huff factor \( h \) \([18]\), which accounts for bound-state corrections arising from Coulomb and relativistic effects. We neglect the Huff factor in the \( p\mu \) system in the following equations, because it is calculated to reduce \( \lambda_{\mu} \) by only 26 ppm \([29,30]\); in the final evaluation of \( \Lambda_{S} \), Eq. \((17)\), we explicitly include this reduction. For the argon system we use \( h = 0.985 \pm 0.003 \), based on extended-model calculations \([19,20]\) that include a more accurate treatment of finite nuclear size effects.

### B. Kinetic equations

The kinetics scheme in Fig. 1 corresponds to a system of coupled linear differential equations for the time-dependent populations \( n_{pp\mu}(t), n_{om}(t), n_{pm}(t), \) and \( n_{Z\mu}(t) \) of the \( p\mu, pp\mu, \) and \( Z\mu \) states. It is convenient to first define the total muon disappearance rate from each state:

\[
\Gamma_{pp\mu} \equiv \lambda_{\mu} + \Lambda_{pZ} + \Lambda_{S} + \Lambda_{pp\mu},
\]

\[
\Gamma_{om} \equiv \lambda_{\mu} + \lambda_{\text{op}} + \Lambda_{OM},
\]

\[
\Gamma_{pm} \equiv \lambda_{\mu} + \Lambda_{PM},
\]

\[
\Gamma_{Z\mu} \equiv h \lambda_{\mu} + \Lambda_{Z}.
\]

These rates are also the eigenvalues of the system. The populations of the muonic states are then described by the
following differential equations:

\[
\begin{align*}
n'_{p\mu}(t) &= -\Gamma_{p\mu} n_{p\mu}(t), \\
n_{om}(t) &= \Lambda_{om} n_{p\mu}(t) - \Gamma_{om} n_{om}(t), \\
n'_{pm}(t) &= \Lambda_{pm} n_{p\mu}(t) + \lambda_{om} n_{om}(t) - \Gamma_{pm} n_{pm}(t), \\
n'_{c\mu}(t) &= \Lambda_{pZ} n_{p\mu}(t) - \Gamma_{Z\mu} n_{Z\mu}(t).
\end{align*}
\]

The initial conditions at \( t = 0 \) are \( n_{p\mu}(0) = f, n_{Z\mu}(0) = f, \) and \( n_{om}(0) = n_{pm}(0) = 0 \). If the rates are time independent, Eq. (7) defines a system of differential equations with constant coefficients that has straightforward but lengthy analytical solutions \( n_i(t) \) (given in the Appendix). Formally they can be written in terms of the eigenvalues \( \Gamma_a \) [see Eqs. (6)]:

\[
n_i(t) = \sum_a \xi_a e^{-\Gamma_a t},
\]

where \( i, \alpha \in (p\mu, \text{om}, pm, Z\mu) \). Usually this is a good approximation, but as is explained later there are cases where epithermal \( p\mu \) atoms are depopulated at energy-dependent rates in the period before they have fully thermalized (cf. Refs. [31,32]). Of particular relevance for the present work, a muonic x-ray measurement [17] observed that the muon transfer rate \( \lambda_{p\mu}(t) \) increased until it reached its constant value for the thermalized atom. Because \( \lambda_{pZ}(t) \) is time dependent, Eq. (7) must be numerically integrated.

From the muonic state populations \( n_i(t) \) we can derive the time distributions of various experimentally observed final-state muon-disappearance products. The distribution of decay electrons is given by

\[
N_e(t) = \lambda_\mu \{ \epsilon_{\mu} [n_{p\mu}(t) + n_{pm}(t)] + \epsilon_{\mu} n_{Z\mu}(t) \},
\]

where \( \epsilon_{\mu} \) is the detection efficiency for electrons produced by muon decay from the hydrogen-bound states. Depending on the experimental setup, the detection efficiency \( \epsilon'_{\mu} \) in higher-\( Z \) atoms can be different because the electron energy spectrum deviates from a pure Michel spectrum due to Coulomb effects [19,20].

The distribution of muon capture products (i.e. recoil nuclei or neutrons) versus time is

\[
N_i(t) = \epsilon_i \{ \Lambda_\mu n_{p\mu}(t) + \Lambda_{om} n_{om}(t) + \Lambda_{pm} n_{pm}(t) \} + \epsilon'_{\mu} \Lambda_{Z\mu} n_{Z\mu}(t).
\]

Here \( \epsilon_i \) and \( \epsilon'_{\mu} \) account for the different efficiencies in detecting reaction products from capture on protons versus capture on nuclei with atomic number \( Z \).

The time distribution of x rays from muon transfer is

\[
N_x(t) = \epsilon_x P_x \Lambda_{pZ} n_{p\mu}(t),
\]

where \( P_x \) is the probability for x-ray emission per transfer and \( \epsilon_x \) is the x-ray detection efficiency. The observables in Eqs. (9)–(11) provide the primary tools for experimentalists in disentangling the rich physics of muon-induced processes in hydrogen.

C. Present experimental knowledge of the molecular formation rate \( \lambda_{pp\mu} \)

The basic experimental technique for measuring the molecular formation rate \( \lambda_{pp\mu} \) is to introduce an impurity to the pure hydrogen target. Though it might seem counterintuitive, adding this complication is helpful because it opens a competing channel to molecular \( pp\mu \) formation. Because muon transfer to the impurity only proceeds from the \( p\mu \) atom, the \( Z\mu \) population follows the time evolution of the \( p\mu \) population that feeds it and the electron distribution described in Eq. (9) depends mainly on \( \lambda_{pp\mu}, \lambda_{pZ} \), and \( \lambda_\mu \). By adding the proper amount of a well-chosen impurity, the terms in Eq. (9) will differ in their time dependencies and relative sizes such that individual rates can be disentangled via a fit to the observed electron time spectrum.

An early measurement of \( \lambda_{pp\mu} \) used an LH\(_2\) target with deuterium admixtures [33]. In this case, muons transfer from \( p\mu \) to \( d\mu \) and deuterium essentially plays the role of the impurity \( Z \) in Fig. 1 and Eq. (9). The formation of \( d\mu \) atoms can lead to muon-catalyzed fusions which emit \( \gamma \)s. Observation of the \( \gamma \) yields for various deuterium concentrations thus enabled a determination of \( \lambda_{pp\mu} \).

Other experiments employed a similar strategy. Conforto et al. [34] measured the muonic x rays emitted following transfer to Ne. Bystritsky et al. [35] simultaneously observed the time distribution of \( \mu Xe \) deexcitation x rays and muon decay electrons. Conforto et al. [34] determined \( \Gamma_{p\mu} \), while Bystritsky et al. [35] enabled independent extraction of \( \lambda_{pp\mu} \) and \( \lambda_{pZ} \) at a single impurity concentration.

The most recent experiment [36] used a very different experimental setup consisting of a layer of solid hydrogen with various tritium admixtures. Fusion products were observed, and muon transfer to tritium changed the disappearance rate of the \( p\mu \) state according to the first of Eqs. (7). Conceptually the experiment was therefore quite similar to Ref. [33].

Figure 2 plots the relevant experimental and theoretical determinations of \( \lambda_{pp\mu} \), including that presented in this paper. The experimental data are not completely consistent. The higher \( \lambda_{pp\mu} \) value measured in the solid-target experiment could originate from comparatively slower thermalization of the \( p\mu \) atoms via elastic collisions with the solid hydrogen lattice [32]. Review [2] excluded the solid-hydrogen result to obtain the experimental world average \( \lambda_{pp\mu} = (2.3 \pm 0.2) \times 10^6 \text{ s}^{-1} \), where the uncertainty has been inflated to account for the inconsistencies among the contributing measurements.

D. Impact of molecular effects on muon capture experiments

Muon capture experiments determine \( \Lambda_\mu \) either by measuring the rate of neutron emission according to Eq. (10) (“neutron method”) or by inferring the muon disappearance rate in hydrogen, \( \lambda_\mu \), from the time distribution of electrons, Eq. (9) (“lifetime method”). While the neutron method does not require high statistics, its precision is fundamentally limited by the fact that the neutron detection efficiency \( \epsilon_c \) must be known to a level that is difficult to achieve in practice. Conversely, the lifetime method requires high statistics but absolute detection efficiency is not a factor. The basic idea of
The lifetime method can be illustrated by considering the ideal case in which only the $p\mu$ state is populated. In that case the electron time distribution Eq. (7) simplifies to

$$N_e(t) \propto e^{-(\lambda_++\lambda_+)t} = e^{-\lambda t}$$

(12)

and $\Lambda_S$ can be determined from the difference $\lambda_- - \lambda_+$. In reality, experiments must always account for effects arising from the existence of muonic molecules. The lifetime method was pioneered by an experiment at Saclay [37] which used an LH$_2$ target ($\phi = 1$); the full kinetics of Eq. (7) therefore needed to be considered, and this led to significant uncertainty in the interpretation of the experiment’s results. The MuCap experiment [4] used a low-density hydrogen target ([33], Conforto et al. [34]), the green cross denotes a solid-target measurement (Mulhauser et al. [36]), and the blue circles denote two measurements in a gaseous hydrogen environment (Bystritsky et al. [35] and this paper). The shaded region corresponds to an updated world average of the experimental results, excluding the outlying solid-target data point.

Figure 3 shows the time distributions of $p\mu$ and $pp\mu$ populations in the hydrogen targets used in the MuCap [4] and Saclay [37] experiments. At the lower density used in the MuCap experiment, muons remain predominantly in the singlet $p\mu$ state over the course of the typical measurement period of 15 $\mu$s. There is nevertheless non-negligible formation of ($pp\mu$)$_{om}$ molecules, and therefore good knowledge of the rate $\lambda_{pp\mu}$ of the process is necessary for correct interpretation of the experiment. In contrast, in the LH$_2$ target used in the Saclay experiment the muon quickly populates the ($pp\mu$)$_{om}$ state, within 1 $\mu$s, and the subsequent depopulation of the ($pp\mu$)$_{om}$ state to the ($pp\mu$)$_{pm}$ state at rate $\lambda_{op}$ is the crucial element to interpreting the experiment.

### III. Experimental Method

#### A. MuCap apparatus

The MuCap detector (Fig. 4) is described here only in brief; greater detail is available in Refs. [4,10,38,39]. The experiment was located at the $\pi E 3$ secondary muon beamline of the 590-MeV proton cyclotron at the Paul Scherrer Institute. Low-energy muons (34 MeV/$c$) passed through a scintillator counter ($\mu$SC) and a wire-chamber plane ($\mu$PC) before coming to a stop inside a 10-bar hydrogen time projection chamber (TPC).

The $\mu$SC provided the start signal for the muon lifetime measurement, and the $\mu$SC and $\mu$PC together provided efficient pileup rejection which enabled selection of events in which only a single muon was present in the TPC. The TPC [40,41] provided tracking of incoming muons and clear identification of each muon’s stopping location by detecting the large peak in energy deposition at the end of the muon’s
B. Measurement of $\lambda_{pp\mu}$ in argon-doped hydrogen

For this measurement we introduced argon to the otherwise ultrapure hydrogen gas, which was at density $\phi = 0.0115 \pm 0.0001$. The atomic concentration of argon was $c_{Ar} = 19.6 \pm 1.1$ ppm, as measured both volumetrically during the initial filling and by gas chromatography at the end of the measurement. The two concentration measurements were consistent, but we have conservatively expanded the uncertainty to cover the uncertainties of both. The gas density was derived from the temperature and pressure, which were continuously monitored.

In principle, two time distributions, $N_c(t)$ [Eq. (10)], and $N_e(t)$ [Eq. (9)], are experimentally observable. The former can be measured either by using the TPC to detect nuclear recoil signals from $\mu + Ar \rightarrow Cl^+ + \nu$ capture events or by using liquid scintillators to detect neutrons emitted by the excited final-state nucleus. There are two disadvantages to measuring $N_e(t)$. First, the spectrum determines $\Gamma_{pp\mu}$ [cf. $n_{Z\mu}(t)$ in Eqs. (A1)] and therefore only the sum of the two transfer rates $\Lambda_{pp\mu}$ and $\Lambda_{pAr}$, not the individual rates themselves, and $\Lambda_{pAr}$ is not known with sufficient precision to enable $\Lambda_{pp\mu}$ to be extracted independently. Second, there are significant systematic uncertainties relating to spatial pileup of TPC signals from the stopping muon and the capture recoil and to uncertainties in the neutron time of flight.

The MuCap experiment was designed to detect decay electrons, so we used a high-statistics sample of $N_e(t)$ to extract $\lambda_{pp\mu}$. If a muon decays it cannot undergo nuclear capture, eliminating the possibility of distortions in muon stop identification due to additional energy deposit from capture recoils. Consequently, the analysis and systematic uncertainties were very similar to those developed for the earlier lifetime experiment measuring $\Lambda_\alpha$ [4].

The decay-electron analysis works as follows. With the judicious choice of argon concentration $c_{Ar} = O(20$ ppm), the disappearance rates $\Gamma_{pp\mu}$ and $\Gamma_{pAr}$ in Eq. (6) are sufficiently different to allow them to be unambiguously extracted from a fit to the corresponding decay-electron time spectrum. The argon capture rate $\lambda_{Ar}$ [14] is 3 times higher than the muon decay rate and therefore transferred muons disappear quickly. Under our conditions, the contributions of $\Lambda_{pp\mu}$ and $\Lambda_{pAr}$ to the total ($p\mu$)S disappearance rate $\Gamma_{pp\mu}$ were 4% and 8%, respectively. As above, the eigenvalue $\Gamma_{pp\mu}$ alone would only determine the sum of two unknowns, $\Lambda_{pp\mu}$ and $\Lambda_{pAr}$. However, both rates enter into the coefficients $c_i^k$ in Eq. (8) in independent combinations, as can be seen from the full solutions in the Appendix. A combined fit can therefore simultaneously determine $\Lambda_{pp\mu}$, $\Lambda_{pAr}$, and, as a byproduct, $\Lambda_{Ar}$ without any need for absolute normalization. To address concerns about the uniqueness and stability of this multiparameter fit to a single distribution, we performed extensive pseudodata Monte Carlo studies of the full kinetics equations; good convergence was observed.

IV. ANALYSIS AND RESULTS

A. Data analysis

A total of $7.2 \times 10^8$ fully reconstructed muon decay events were used in the present analysis. These events were selected via application of our standard cuts, described in Ref. [4]. Each event was required to involve a pileup-free muon stop in the TPC fiducial volume, $\Delta x \times \Delta y \times \Delta z = 10.4 \times 8.0 \times 20.4 \text{ cm}^3$. The decay-electron trajectories were reconstructed from spatial and temporal coincidences among the two cylindrical wire chambers and the two layers of plastic scintillators. Once the set of good events had been selected, the time differences between the fast signals of the electron scintillator eSC and the muon beam scintillator $\mu$SC were histogrammed and the resulting decay time spectrum was fitted with the function

$$N(t) = \lambda [n_{pm}(t) + n_{pm}(t) + n_{om}(t) + \epsilon \ h \ n_{Z\mu}(t)] + B$$

using the MINOS package. This fit function is identical to Eq. (9) apart from the introduction of the flat background term $B$. The relative efficiency $\epsilon$ is defined as $\epsilon \equiv \epsilon_c/\epsilon_e$.

To accommodate the time dependence of $\lambda_{pAr}$ in a nearly model-independent way, this rate was parametrized in the form

$$\lambda_{pAr}^\text{fit}(t) = \lambda_{pAr}(1 - \alpha e^{-\beta t}),$$

where $\alpha$ and $\beta$ were extracted from Fig. 1 in Ref. [17]. The parameter $\beta$ characterizes $p\mu$ thermalization and was scaled down by 1.5 from the value in Ref. [17], as that experiment used a 15-bar target whereas MuCap used a 10-bar target. The scaling of $\alpha$ with pressure depends on the initial population of hot $p\mu$ atoms after the muonic cascade, which, according to theory [42], should increase by ~10% with a pressure increase from 10 to 15 bar. We did not change the value of $\alpha$ extracted from Ref. [17], but we assigned it a conservative 50% uncertainty. The uncertainties in $\alpha$ and $\beta$ in Table II were used to extract the effect on the fitted rates. The resulting systematic

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>$c_{Ar}$</td>
<td>$19.6 \pm 1.1$ ppm</td>
</tr>
<tr>
<td>$\phi$</td>
<td>$0.0115 \pm 0.0001$</td>
</tr>
<tr>
<td>$f$</td>
<td>$5 \pm 1 \times 10^{-4}$</td>
</tr>
<tr>
<td>$\epsilon$</td>
<td>$0.996 \pm 0.003$</td>
</tr>
<tr>
<td>$c_{O}$</td>
<td>$57 \pm 57$ ppb</td>
</tr>
<tr>
<td>$c_{N}$</td>
<td>$115 \pm 115$ ppb</td>
</tr>
<tr>
<td>$\alpha$</td>
<td>$0.25 \pm 0.12$</td>
</tr>
<tr>
<td>$\beta$</td>
<td>$1.0 \pm 0.2 \times 10^7$ s$^{-1}$</td>
</tr>
</tbody>
</table>
TABLE III. Fit results for $\Lambda_{pp\mu}$, $\Lambda_{pAr}$, and $\Lambda_{Ar}$, as well as their associated systematic corrections ($\Delta$) and uncertainties ($\delta$). The final error on each rate is the quadrature sum of the contributing uncertainties.

<table>
<thead>
<tr>
<th>Fit</th>
<th>$\Lambda_{pp\mu}$ (s$^{-1}$)</th>
<th>$\Lambda_{pAr}$ (s$^{-1}$)</th>
<th>$\Lambda_{Ar}$ (10$^2$ s$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Systematic</td>
<td>$\Delta$</td>
<td>$\delta$</td>
<td>$\Delta$</td>
</tr>
<tr>
<td>Timing calibration</td>
<td>39</td>
<td>13</td>
<td>35</td>
</tr>
<tr>
<td>Efficiency $\epsilon$</td>
<td>37</td>
<td>13</td>
<td>34</td>
</tr>
<tr>
<td>Huff factor $h$</td>
<td>45</td>
<td>15</td>
<td>27</td>
</tr>
<tr>
<td>$f$</td>
<td>46</td>
<td>12</td>
<td>39</td>
</tr>
<tr>
<td>$\Lambda_{N},\Lambda_{OM,PM}$</td>
<td>13</td>
<td>16</td>
<td>5</td>
</tr>
<tr>
<td>$\lambda_{op}$</td>
<td>31</td>
<td>3</td>
<td>2</td>
</tr>
<tr>
<td>$\Lambda_{pf}$</td>
<td>9</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Epithermal</td>
<td>329</td>
<td>85</td>
<td>278</td>
</tr>
<tr>
<td>H$_2$O and N$_2$</td>
<td>116</td>
<td>116</td>
<td>$-15$</td>
</tr>
<tr>
<td>Final result</td>
<td>23 112</td>
<td>$\pm$ 741</td>
<td>43 784</td>
</tr>
</tbody>
</table>

Uncertainties are listed under “Epithermal” in Table III. The final fit method used numerical integration with the values listed in Table II. The analytical solution (A2) was used for cross checks.

The fitting procedure using Eq. (13) requires a timing calibration to assert that the muon arrival time is at $t = 0$. For that, the rising edge of the histogrammed differences of the $\mu$SC and the 16 eSC subdetectors were fitted individually. This determined timing calibration offsets for each eSC detector with a precision of 2 ns. The 16 offsets were then applied to their corresponding spectrum before the sum of all time distributions was fit with Eq. (13).

The fit was performed over the range [0.12 $\mu$s, 20 $\mu$s]. Five quantities were treated as free parameters: $\Lambda_{pp\mu}$, $\Lambda_{pAr}$, $\Lambda_{Ar}$, the normalization $A$, and the background term $B$. All other parameters were fixed in the fit to the values in Table I and, for experiment-specific parameters, according to the values given in Table II. The initial $\mu$Ar formation fraction $f = (5 \pm 1) \times 10^{-4}$ is the sum of two components, $f_e$ and $f_c$. The atomic capture ratio for argon relative to hydrogen has been measured to be $f_c = (9.5 \pm 1.0)_{\text{arg}} = (1.87 \pm 0.20) \times 10^{-4}$ [43]. An additional initial population $f_e = (1.66 \pm 0.34)$, from excited-$\mu$-state transfer has been observed in a target at 15-bar pressure [43]. We account for this by using $f_e = (3.1 \pm 0.9) \times 10^{-4}$, in which the uncertainty has been conservatively enlarged to accommodate the possibility of a pressure dependence.

The energy spectra of decay electrons emitted from $p\mu$ and $Ar\mu$ atoms are different, which leads to a difference in the corresponding detection efficiencies. We used the energy spectrum calculated in Ref. [44] and folded it together with the energy-dependent detector efficiency obtained from a full GEANT4 simulation. The resulting relative efficiency, $\epsilon = 0.996 \pm 0.003$, shows that the thin layers of the MuCap electron detectors are not very sensitive to spectral differences at higher energies.

After the fit, small corrections were applied to the fitted rates to account for the presence of the chemical impurities oxygen and nitrogen, with atomic concentrations $c_O$ and $c_N$, respectively. This procedure is discussed in the next section.

B. Results and systematic uncertainties

The fit to the data is plotted in Fig. 5. Figure 5(a) shows the decay-electron time spectrum alongside the time distributions of the parent muon populations $n_{p\mu}(t)$, $n_{Ar\mu}(t)$, $n_{on}(t)$, and $n_{pm}(t)$ determined by the fit. Figure 5(b) displays the residuals, i.e., the differences between the data and the fit function normalized by the uncertainty of each data point. The good agreement between the data and the fit function is demonstrated by the reduced $\chi^2$/DOF = 0.983 ± 0.064.

Table III presents the fit results for the three rates $\Lambda_{pp\mu}$, $\Lambda_{pAr}$, and $\Lambda_{Ar}$. The table also lists systematic corrections $\Delta$ and the systematic uncertainties $\delta$ resulting from a ±1σ variation of the fixed parameters listed in Tables I and II.

The fit did not explicitly model effects from the accumulation of nitrogen and oxygen in the hydrogen due to outgassing from the TPC vessel. Instead, a correction $\Delta$ was applied to the fitted values of both $\Lambda_{pp\mu}$ and $\Lambda_{pAr}$. During its main run MuCap achieved hydrogen chemical purity levels of better than 10 ppb, but during the argon-doped measurement the TPC was disconnected from the hydrogen circulation and purification system [45]. After 6 days, atomic concentrations of $c_O = 115$ ppb of oxygen (in the form of water vapor) and $c_N = 230$ ppb of nitrogen were observed using a humidity sensor and gas chromatography, respectively. Due to the higher muon transfer and capture rates for oxygen compared to nitrogen,
The pseudodata were then fitted with Eq. (13) to extract the parameters in the fit to the decay-electron time spectrum. The normalized correlations among the five free fit parameters are presented in Table IV. These correlations are statistically allowed and cover all possible accumulation scenarios. The impurity-related corrections to \(\lambda_{pp\mu}\) and \(\lambda_{pAr}\) were determined to be \(\Delta_{pp\mu} = 116 \pm 116 \, \text{s}^{-1}\) and \(\Delta_{pAr} = -15 \pm 15 \, \text{s}^{-1}\).

The final results for the fitted rates after applying the impurity-related corrections and summing all systematic uncertainties (Table III) are

\[
\begin{align*}
\lambda_{pp\mu} &= 2.311 \pm 0.074 \times 10^4 \, \text{s}^{-1}, \\
\lambda_{pAr} &= 4.378 \pm 0.018 \times 10^4 \, \text{s}^{-1}, \\
\lambda_{Ar} &= 1.302 \pm 0.032 \times 10^6 \, \text{s}^{-1}.
\end{align*}
\]

From these one can deduce the normalized rates

\[
\begin{align*}
\lambda_{pp\mu} &= 2.01 \pm 0.07 \times 10^6 \, \text{s}^{-1}, \\
\lambda_{pAr} &= 1.94 \pm 0.11 \times 10^{11} \, \text{s}^{-1},
\end{align*}
\]

using Eqs. (3) and (5), respectively.

The normalized correlations among the five free fit parameters are presented in Table IV. These correlations are incorporated into the uncertainties on the final results.

Our result for \(\lambda_{pp\mu}\) is about 1\(\sigma\) larger than the value we obtained in Ref. [4] due to the more refined analysis in this paper and the correction of a numerical error in the fitting code. As regards the transfer rate to argon \(\lambda_{pAr}\), there is a wide spread of experimental results obtained with different methods and target conditions, clustered around \(1.4 \times 10^{11} \, \text{s}^{-1}\), \(3.6 \times 10^{11} \, \text{s}^{-1}\), and \(9 \times 10^{11} \, \text{s}^{-1}\), as discussed in Ref. [17]. Our value \(\lambda_{pAr} = 1.94 \pm 0.11 \times 10^{11} \, \text{s}^{-1}\) is close to the most recently published value, \(1.63 \pm 0.09 \times 10^{11} \, \text{s}^{-1}[17]\), albeit 2.2\(\sigma\) higher. Note that the uncertainty in the argon concentration only enters into the extraction of the normalized rate \(\lambda_{pAr}\), while in the fit to determine \(\lambda_{pp\mu}\) effective rates are being used that are independent of \(c_\text{Ar}\). Our result for the muon’s nuclear capture rate on argon, \(\lambda_{Ar}\), agrees well with the values in the literature: \(1.20 \pm 0.08 \times 10^6 \, \text{s}^{-1}\) [46] and \(1.41 \pm 0.11 \times 10^6 \, \text{s}^{-1}\) [47].

C. Consistency checks

The fit start time was varied to check for any distortions or physical effects not accounted for by the fit function. Figure 6

![Figure 6](image_url)
shows the progressions of the fitted rates as the fit start time was increased in steps from its standard value of 0.12 μs. The red lines denote the ±1σ variation allowed because of the set-subset statistics involved in this procedure. Each rate is statistically self-consistent across the fit start time scan.

In the fit to the data using Eqs. (9) and (13), the capture rate \( \Lambda_S \) is required as an input to extract \( \lambda_{pp\mu} \), while the latter is itself used in the determination of \( \Lambda_S \). This interdependence is not a problem because all fitted rates in Eq. (15) depend only very weakly on the hydrogen capture rates, as quantified in Table III. We explicitly iterated the procedure (obtaining fit results with \( \Lambda_S \) as input, using the results to correct \( \Lambda_S \), repeating with the adjusted \( \Lambda_S \)) to arrive at a stable, self-consistent solution, and we found that the results for \( \lambda_{pp\mu} \), \( \Lambda_{p\mu\mu} \), and \( \Lambda_{Ar} \) were changed by less than one-tenth of their uncertainties.

Last, the reproducibility of the fit was tested by generating 10⁴ pseudodata histograms using the final fit parameters in Eqs. (15) and fitting each pseudoeperiment in the same manner as the real data. The fits consistently yielded the input values, and the simulated data reproduced the same fit uncertainties listed in Table III.

V. RELEVANCE TO THE INTERPRETATION OF THE MUCAP EXPERIMENT

In Sec. II the influence of the molecular rates \( \lambda_{pp\mu} \) and \( \lambda_{op} \) on muon kinetics in hydrogen was described. The MuCap experiment measured the effective muon disappearance rate \( \lambda_\perp \) in low-density ultrapure hydrogen by fitting the observed decay electron time distribution with a three-parameter function, \( f(t) = A e^{-\lambda_\perp t} + B \). Taking \( pp\mu \) formation into account, the disappearance rate can be expressed as

\[
\lambda_\perp = \lambda_\parallel + \Delta \lambda_{pp\mu} + \Lambda_\perp + \Delta \Lambda_{pp\mu}. \quad (17)
\]

Here \( \Delta \lambda_{pp\mu} = -12.3 \text{ s}^{-1} \) is a calculable bound-state modification to the muon decay rate in the \( p\mu \) system \( \text{(29, 30)} \), and \( \Delta \Lambda_{pp\mu} \) is a modification to \( \Lambda_\perp \) accounting for the small population of muonic molecules and the fact they have unique capture rates. In the following we summarize the derivation of \( \Delta \Lambda_{pp\mu} \), based on our improved measurement of \( \lambda_{pp\mu} \) at conditions nearly identical to those of the main MuCap experiment.

The derivation is based on high-statistics simulations of the full kinetics described by Eqs. (11). Because the MuCap measurement of \( \Lambda_\perp \) was performed using pure hydrogen gas, for the simulations the Z channel was used to model the small amounts (few ppb) of oxygen and nitrogen impurities that were observed to have outgassed from the hydrogen vessel’s walls. The relevant input parameters for the simulation were those in Tables I and II. An accidental background was added to make the signal-to-background level commensurate with that in the MuCap data. Time distributions of \( 10^{12} \) decay electrons were generated for two different cases: \( \lambda_{pp\mu} = 0 \) and \( \lambda_{pp\mu}^\text{MuCap} \). The previous MuCap analysis \( \text{(4)} \) was performed using the preliminary value \( \lambda_{pp\mu}^\text{MuCap} = 1.94 \pm 0.06 \times 10^6 \text{ s}^{-1} \); here we update the analysis using our new result in Eq. (21).

To determine the effect on the MuCap result for \( \Lambda_\perp \), we fit the simulated time distributions with the same three-parameter function used to fit the data. The relevant correction is then obtained via

\[
\Delta \Lambda_{pp\mu} = \lambda_\perp (\lambda_{pp\mu}^\text{MuCap}) - \lambda_\perp (\lambda_{pp\mu} = 0), \quad (18)
\]

where the \( \lambda_\perp \) values are obtained from fits to the two simulated data sets generated using different \( \lambda_{pp\mu} \) values. The uncertainty in \( \Delta \Lambda_{pp\mu} \) is estimated in a similar manner, by generating pseudodata while varying the parameters entering the kinetic equations by ±1σ individually. The resulting fit determines the final correction for the MuCap experiment to be \( \Delta \Lambda_{pp\mu} = -18.4 \pm 1.9 \text{ s}^{-1} \), which is smaller than the correction in Ref. [4] by 0.7 s⁻¹. Thus the updated value of \( \lambda_{pp\mu}^\text{MuCap} \) induces a small shift of the singlet \( p\mu \) capture rate measured by MuCap from \( \Lambda_\perp = 714.9 \pm 5.4 \text{stat} \pm 5.1 \text{syst s}^{-1} \) obtained in Ref. [4] to

\[
\Lambda_\perp = 715.6 \pm 5.4 \text{stat} \pm 5.1 \text{syst s}^{-1}. \quad (19)
\]

The value of the pseudoscalar coupling constant, \( g_p = 8.06 \pm 0.55 \) extracted in Ref. [4], is correspondingly changed by -0.045, i.e., by only 8% of its uncertainty.

From our simulations we can determine the dependence of \( \Delta \Lambda_{pp\mu} \) on molecular parameters:

\[
\Delta \Lambda_{pp\mu} = -18.4 \left[ 1 + a (\lambda_{pp\mu} - \lambda_{pp\mu}^\text{MuCap}) + b (\lambda_{op} - \lambda_{op}^0) \right], \quad (20)
\]

where \( \lambda_{op}^0 \) is given in Table I, \( a = 4.7 \times 10^{-7} \), and \( b = 2.9 \times 10^{-6} \). Using the new measurement presented in this paper, the total uncertainty in the MuCap capture rate \( \Lambda_\perp \) due to \( pp\mu \) formation is less than 2 s⁻¹ and is dominated by \( \lambda_{op} \), while \( \lambda_{pp\mu} \) contributes only 0.6 s⁻¹.

VI. SUMMARY

The time spectrum of electrons emitted by the decay of muons stopped in argon-doped hydrogen was measured with the MuCap detector, for the purpose of determining the formation rate \( \lambda_{pp\mu} \) of \( pp\mu \) muonic molecules. The TPC enabled selection of muons that stopped in the hydrogen, away from high-Z materials, and the electron tracker provided 3π solid-angle coverage and enabled vertex matching with muon stops. We developed a detailed physics model to describe the time evolution of the atomic and molecular muonic states contributing to the decay-electron spectrum, taking into account the energy dependence of the muon transfer rate \( \lambda_{pAr} \) from hydrogen to argon. We extracted \( \lambda_{pp\mu} \), \( \lambda_{pAr} \), and the muon capture rate in argon, \( \Lambda_{Ar} \), from a single fit to the decay-electron time spectrum. Our results for \( \lambda_{pAr} \) and \( \Lambda_{Ar} \) agree with those from previous dedicated experiments. Our result for the \( pp\mu \) formation rate,

\[
\lambda_{pp\mu}^\text{MuCap} = 2.01 \pm 0.06 \text{stat} \pm 0.03 \text{syst} \times 10^6 \text{ s}^{-1}, \quad (21)
\]

is 2.5 times more precise than previous measurements, which were performed under a variety of different experimental conditions and whose results disagreed beyond their uncertainties. To obtain a new world average we used the procedures for averaging and inflating uncertainties advocated by the
Particle Data Group [11]; we have included only the gas- and liquid-target experiments [33–35], choosing to omit the lone solid-target experiment [36] because of possible solid-state effects that are not well understood. The updated experiment world average then becomes

$$\lambda_{pp\mu}^{avg} = 2.10 \pm 0.11 \times 10^6 \text{ s}^{-1}. \quad (22)$$

The rate $\lambda_{pp\mu}$ was a necessary input to the MuCap experiment’s recent precision determination of the nuclear capture rate on the proton, $\Lambda_S$ [4]. The MuCap experiment was designed so that the majority of muons underwent capture in muonic $p\mu$ atoms, and formation of $pp\mu$ molecules changed the observed capture rate by only 2.5%. However, given the inconsistency between existing $\lambda_{pp\mu}$ results it was difficult to confidently estimate the uncertainty on the correction to $\Lambda_S$ for $pp\mu$ effects. Our new result for $\lambda_{pp\mu}$, obtained at the same hydrogen density and temperature as in the main MuCap experiment, leads to a well-defined correction to $\Lambda_S$, and the corresponding contribution to the total error is now minor. The value of $\lambda_{pp\mu}$ presented here differs only slightly from the value used in Ref. [4], and consequently the updated values for $\Lambda_S$ and the pseudoscalar coupling $g_p$ agree to better than 0.1σ with the values in that publication.

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APPENDIX: SOLUTIONS TO THE MUON KINETICS EQUATIONS

The differential equations in Eq. (7) can be solved by determining the eigenvalues and eigenvectors of the system. The time-dependent populations of the four muonic states are given by

$$n_{pm}(t) = (1 - f)e^{-\Gamma_{pm}t},$$

$$n_{om}(t) = (1 - f)\frac{\Lambda_{om}}{\Gamma_{om} - \Gamma_{pm}}(e^{-\Gamma_{pm}t} - e^{-\Gamma_{om}t}),$$

$$n_{pp}(t) = \frac{1 - f}{\Gamma_{pm} - \Gamma_{pp}} \left[ \frac{\Lambda_{om} \Lambda_{pm} \Lambda_{pp}}{\Gamma_{pm} - \Gamma_{om}} \left( e^{-\Gamma_{pm}t} - e^{-\Gamma_{om}t} \right) + \frac{\Gamma_{pp} \Lambda_{pp} \Lambda_{om} \Lambda_{pm}}{\Gamma_{pm} - \Gamma_{pp}} \right],$$

$$n_{Z\mu}(t) = (1 - f)\frac{\Lambda_{Z\mu} \Lambda_{pp}}{\Gamma_{Z\mu} - \Gamma_{pp}}(e^{-\Gamma_{pm}t} - e^{-\Gamma_{Z\mu}t}) + f e^{-\Gamma_{Z\mu}t}.$$ (A1)

One simplistic but heuristic approximation is to neglect the small parameters $\Lambda_S$ and $\Lambda_{pp\mu}$ in the disappearance rates in Eqs. (6) and the initial $Z\mu$ population $f$ and to assume that $\Gamma_{Z\mu}$ is large compared to all other eigenvalues in Eq. (6). In this limit the observable electron distribution, Eq. (9), attains the simple form

$$N_e(t) \propto e^{-\lambda_{pp\mu}t} \left[ 1 + \frac{\Lambda_{pp\mu}}{\Lambda_{pp\mu} + \Lambda_{pp\mu}^2} \right], \quad (A2)$$

which elucidates our strategy of determining $\Lambda_{pp\mu}$ in a single fit.