On the Enhanced Cosmic-Ray Ionization Rate in the Diffuse Cloud toward ζ Persei

Gargi Shaw
University of Kentucky

Gary J. Ferland
University of Kentucky, gary@uky.edu

R. Srianand
IUCAA, India

N. P. Abel
University of Kentucky

P. A. M. van Hoof
Royal Observatory of Belgium, Belgium

See next page for additional authors

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Authors
Gargi Shaw, Gary J. Ferland, R. Srianand, N. P. Abel, P. A. M. van Hoof, and P. C. Stancil

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ON THE ENHANCED COSMIC-RAY IONIZATION RATE IN THE DIFFUSE CLOUD TOWARD ζ PERSEI

GARGI SHAW,1 G. J. FERLAND,1 R. SRIANAND,2 N. P. ABEL,1 P. A. M. VAN HOOF,3 AND P. C. STANCIL4

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ABSTRACT

The spatial distribution of the cosmic-ray flux is important in understanding the interstellar medium (ISM) of the Galaxy. This distribution can be analyzed by studying different molecular species along different sight lines whose abundances are sensitive to the cosmic-ray ionization rate. Recently several groups have reported an enhanced cosmic-ray ionization rate (ζ = ζCR/ζstandard) in diffuse clouds compared to the standard value, ζstandard(=2.5 × 10⁻¹⁷ s⁻¹), measured toward dense molecular clouds. In an earlier work we reported an enhanced ζCR = 20 toward HD 185418. McCall et al. have reported ζCR = 48 toward ζ Persei based on the observed abundance of H₃⁺, while Le Petit et al. found ζCR ≈ 10 to be consistent with their models for this same sight line. Here we revisit ζ Persei and perform a detailed calculation using a self-consistent treatment of the hydrogen chemistry, grain physics, energy and ionization balance, and excitation physics. We show that the value of ζCR deduced from the H₃⁺ column density, N(H₃⁺), in the diffuse region of the sight line depends strongly on the properties of the grains because they remove free electrons and change the hydrogen chemistry. The observations are largely consistent with ζCR ≈ 40, with several diagnostics indicating higher values. This underscores the importance of a full treatment of grain physics in studies of interstellar chemistry.

Subject headings: cosmic rays — ISM: abundances — ISM: clouds — ISM: individual (ζ Persei) — ISM: structure

1. INTRODUCTION

ζ Persei is a type B star (Snow 1977; Perryman et al. 1997) situated at a distance of 300 pc (R.A.: 03°58′57.93″, decl.: +35°47′28″, J2000.0). The diffuse cloud toward ζ Persei is important because of the measurement of a large H₂ column density, N(H₂). H₃⁺ plays an important role in the ion-molecule chemistry in the interstellar medium (ISM). It is created mainly via the gas-phase reaction H₂ + H₂ → H₃⁺ + H once H₂ is produced by cosmic-ray ionization of H₂. Therefore, a large N(H₃⁺) value suggests the presence of an enhanced cosmic-ray ionization rate along this sight line (McCall et al. 2003, hereafter MC03). A nice chronological discussion of the problem can be found in Le Petit et al. (2004; hereafter LP04).

Several groups have proposed that an enhanced cosmic-ray ionization rate, ζCR(ζ/ζstandard, where ζstandard is the standard rate), may vary across the Galaxy (e.g., Federman et al. 1996). MC03 have derived an enhanced ζCR = 48 toward ζ Persei by combining an experimental determination of the dissociative recombination rate of H₃⁺(kₜ) with the observed N(H₃⁺) and assuming a hydrogen number density of 250 cm⁻³. On the other hand, LP04 modeled the diffuse cloud as a two-phase plane-parallel slab that is illuminated on one side by twice the galactic background radiation field. They found that ζCR = 10 can reproduce most column densities within a factor of 3, but they required shocks to reproduce the observed N(CH₃) and the populations of excited rotational levels of H₂(J = 2, 3, 4, and 5). Hence they concluded that enhancement of ζCR is less than that required by MC03.

In an earlier paper (Shaw et al. 2006, hereafter S06), we showed that when the particle density is well constrained by fine-structure level populations, the observed ortho-para temperature and high-J excitations of H₂ are very sensitive to ζCR. Further, it was found in S06 that ζCR = 20 reproduces the observed temperature and high-J excitation in the diffuse ISM toward HD 185418 without requiring additional shocked gas.

Performing an independent analysis is one way of verifying the high ζCR requirement toward ζ Persei. In a diffuse cloud, H₃⁺ is mainly destroyed by dissociative recombination with electrons (H₃⁺ + e⁻ → H + H + H or H₂ + H). Assuming steady state, kₜ = kₜ(n(H₃⁺))nₑ / n(H₂),

where n(H₂), n(H₃⁺), and nₑ are densities of H₂, H₃⁺, and e⁻, respectively. For a given kₜ and n(H₂), the relationship between (H₃⁺) and ζCR depends on nₑ. To determine all of the parameters necessary to evaluate ζCR, we perform a detailed calculation, using a self-consistent treatment of the hydrogen chemistry, grain physics, energy and ionization balance, and excitation physics. We find that grains are very important in determining the free electron density, nₑ. The grains along this sight line have a small ratio of total to selective extinction Rₑ, indicating a grain size distribution with an enhanced number of small particles. The small grains are very effective at absorbing free electrons, changing the ionization balance of the cloud. This alters the ζCR deduced from the observed N(H₃⁺), as illustrated below.

2. DESCRIPTION OF THE MODEL AND PARAMETERS

2.1. Chemical and Thermal State

In this work, we used version 07.02 of the spectral simulation code Cloudy, which was last described by Ferland et al. (1998). It performs a self-consistent calculation of the thermal, ionization, and chemical balance of both the gas and dust. The chemical network consists of ~10^5 reactions with 71 species involving H,
He, C, N, O, Si, S, and Cl (Abel et al. 2005). The effects of ion charge transfer on grain surfaces are described in Abel et al. (2005). Comparison with other existing codes (Roellig et al. 2007) finds good agreement in the PDR (photo-dominated region) limit.

Further, we include silicate and graphite size-resolved grains and determine the grain charge and photoelectric heating self-consistently (van Hoof et al. 2004). The calculations use 10 size bins each for silicate and graphite. We adjust the size distribution to reproduce the observed extinction, as discussed below. We do not include polycyclic aromatic hydrocarbons (PAHs) due to a lack of observational evidence along this line of sight.

Our treatment of H₂ is described in Shaw et al. (2005, hereafter S05). The H₂ chemistry network consists of various state-specific formation and destruction processes. The model explicitly includes cosmic-ray excitation to the singlet states and triplet states via secondary electrons generated by cosmic rays (Dalgarno et al. 1999). Here we improve this model by adopting, for the singlet states B, C, B', and D, the Born approximation cross section

\[ \sigma = 4 \pi \frac{m_e^2}{2 \pi a_0} \frac{1}{k^2} \ln \left( \frac{k' + k}{k' - k} \right) \]  

(Liu & Dalgarno 1994) in terms of the dipole-allowed electronic spontaneous emission transition probabilities. The k and k' values are the incoming and outgoing momenta of the electrons, and all other symbols have their usual meanings. We assume an average incoming energy of 20 eV. The excitation rate is estimated by rescaling this cross section in terms of hydrogen ionization cross section (Shemansky et al. 1985) and then rescaling with the cosmic-ray excitation rate of Lyα. The cross section for the dissociative triplet state b is taken from Dalgarno et al. (1999, Fig. 4b), rescaled in terms of the hydrogen ionization cross section (Shemansky et al. 1985), and then multiplied with the cosmic-ray excitation rate of Lyα. These processes excite the electronic states, eventually populating the higher vib-rotational levels of the ground state, thus changing the populations of levels within X.

In addition to the chemistry described in Abel et al. (2005), we now consider condensation of CO, H₂O, and OH onto grain surfaces following the formalism of Hasegawa et al. (1992), Hasegawa & Herbst (1993), and Bergin et al. (1995). For each molecule, condensation is balanced with desorption due to thermal and cosmic-ray evaporation. Grain surface reactions between molecules are not considered. Condensation is important when the gas and dust temperature fall below 20–25 K (Bergin et al. 1995). Removing a molecule from the gas phase will alter many of the reaction rates and will cause the molecular abundances to change from the case of pure gas-phase chemistry. Although included, it was found that molecules do not condensate on grain surfaces in the environments considered here as the grain temperature is higher than 25 K.

Bound-electron Compton scattering is the most efficient photoionization process for photon energies above ~10 keV (Osterbrock & Ferland 2006, Fig 11.5). We assume a cross section that is scaled from the Klein-Nishina electron cross section multiplied by the number of valence electrons. This is included as a general ionization and heating process for all molecules and for molecules condensed onto grain surfaces.

The gas kinetic temperature is determined at each point in the cloud, including all relevant heating and cooling processes as detailed in Abel et al. (2005). In a typical calculation, the temperature will range from ~5000 K, where a significant amount of H⁺ is present, to ~50 K in well-shielded molecular regions. The cloud is assumed to be sufficiently old for the chemical, ionization, and thermal balances to have reached time-steady conditions.

2.2. Focus of this Paper

Gas is present in three distinct phases (ionized, diffuse, and dense) along this sight line (Snow 1977; van Dishoeck & Black 1986; LP04). Most of the column density is in a low-density PDR-like medium containing H₂ and C⁺. Thin layers that are predominantly molecular or ionized are also present. Our emphasis is on using N(H₂) to derive χCR, while simultaneously accounting for other observables that are produced in the same environment where H₂ forms. The abundance of H⁺ is proportional to the H₂ abundance and inversely proportional to n_e, so we concentrate on the diffuse phase where most of the H₂ resides and investigate different processes that influence n_e.

The observed N(C2) and N(CN) are produced in a thin dense molecular phase (Gredel et al. 2002). The total column density of this phase is only N(H₂) > 10¹⁹ cm⁻² (LP04) and should, therefore, give only a small contribution to the total N(H₂). In what follows, we model the diffuse phase and do not attempt to model the conditions in the dense molecular phase. We therefore make no attempt to reproduce the N(C2) and N(CN) column densities and will not further develop the LP04 model of the dense phase.

The presence of a significant column density of Si¹⁺ and S¹⁺ indicates that an H⁺ region is also present along this sight line. As described below, the radiation field includes a hydrogen-ionizing continuum (Black 1987), which produces a thin layer of H⁺ on the surface of the diffuse cloud. Nearby hot stars contribute additional hydrogen-ionizing photons. These ions are attributed to an H⁺ region near the background star as in S06 and no attempt is made to reproduce them here.

2.3. Cloud Geometry and Illuminating Radiation

We consider a plane-parallel geometry with radiation striking both sides. Van Dishoeck & Black (1986) found that this geometry reproduced the physical conditions within typical diffuse interstellar clouds. We also used this geometry to reproduce the properties of a well-studied sight line toward HD 185418 (S06). The Galactic background radiation field given by Black (1987) and the cosmic rays are the only source of photoelectric heating and ionization in the calculations. We interpolate an FUV continuum between the UV and X-ray observations (13.6–54 eV) and our interpolation of Black (1987) is shown in Table 1. We parameterize the intensity of the incident radiation by χ, the ratio of the assumed incident radiation field to the Galactic background. The scale factor multiplies the entire continuum, except for the CMB (cosmic microwave background) which is added separately to the rescaled ISM field. The radiative transport of the full continuum, from radio through X-rays, is done including pumping by the incident continuum, atomic and molecular lines, and background opacities. It also includes line overlap by the hundreds of thousands of electronic transitions that determine the destruction and excitation of H₂ (S05) and other lines.

2.4. Model Parameters

We have adopted R_e = 2.85 (Cardelli et al. 1989) and E(B − V) = 0.32, giving A_v/N(H₂) = 5.7 × 10⁻²² cm², where N(H₂) is the column density of H₂ in all forms. The grain-size distribution was modified to approximate this extinction law. Following Mathis et al. (1977), a power-law size distribution of grains was adopted, dn/da ∝ a⁻³.², where a is the radius of the grain and dn is the number of grains with a radius lying within a and a + da.
The total H$_2$ column density is not modified by turbulence because most of the turbulence but keeping all the other parameters identical to our physical microturbulence but could include macroscopic motion for the H$_2$ lines (Snow 1977; LP04). This may not actually be the case, as the volume is 

The density is determined by matching the observed column densities of C, (van Hoof et al. 2004). We use the size-dependent grain temperatures of the grains are determined self-consistently (graphite with 10 size bins ranging between 10$^2$ cm$^{-3}$ since they are obtained self-consistently. We find a total hydrogen density of n(H$_{\text{tot}}$) = 80 cm$^{-3}$. The Galactic starlight background enhancement $\chi$ is obtained by fitting S$^0$/S$^1$, Mg$^0$/Mg$^+$, Fe$^0$/Fe$^+$, C$^0$/C$^+$, and Ca$^0$/Ca$^+$. We find $\chi$ close to 2.

The resulting parameters are listed in Table 2 and are reasonably close to those found by LP04. This best-fitting model will be described further below.

The remaining free parameter is the cosmic-ray ionization rate enhancement $\chi_{\text{CR}}$. This is obtained by matching to the column densities of CO, H$^0$, and H$_2$ in various J levels, NH, and H$^+$. In additional calculations, all the parameters listed in Table 2 are kept fixed, and $\chi_{\text{CR}}$ is varied to illustrate its effect on the predicted column densities. We find that S$^0$/S$^1$, Mg$^0$/Mg$^+$, Fe$^0$/Fe$^+$, C$^0$/C$^+$, and Ca$^0$/Ca$^+$ have a small dependence on $\chi_{\text{CR}}$ in addition to their dependence on $\chi$. These ionization ratios change by less than 0.2 dex for 1 < $\chi_{\text{CR}}$ < 100, while the ratios change nearly linearly with $\chi$. The 2 $\sigma$ errors in the observed values are larger than 0.2 dex, so we did not use them to constrain $\chi_{\text{CR}}$.

3. RESULTS AND DISCUSSION

3.1. A Model with $\chi_{\text{CR}} = 1$

As a first test, we consider $\chi_{\text{CR}} = 1$ and compute the hydrogen-ionization structure as a function of cloud depth. We assume the density and radiation field derived in the previous section. Figure 1 shows half of the symmetrically ionized face. We stopped our calculation at half of the observed value of N(H$_{\text{tot}}$) = 10$^{21.2}$ cm$^{-2}$ since the cloud is assumed to be symmetric. At the illuminated face the gas is moderately ionized, with n(H$^0$) $\approx$ 35 cm$^{-3}$ and n(H$^+$) $\approx$ 45 cm$^{-3}$. In deeper regions hydrogen is molecular with the H$_2$ fraction f(H$_2$) $\approx$ 0.99. The electron density $n_e$ and electron temperature $T_e$ are displayed as functions of cloud depth in Figure 2. The density $n_e$ varies from 45 to 0.014 cm$^{-3}$ while the temperature ranges from 5370 to 38 K across the cloud. We see that H$^+$ is the main electron donor at the illuminated face. The predicted column densities are listed in the third column of Table 3. This model clearly underpredicts N(H$^0$), N(C$^0$), N(NH), N(CO), N(OH), N(H$^+_2$), and the column densities of the J $\geq$ 2 rotational levels of H$_2$.

3.2. Models with Varying $\chi_{\text{CR}}$

Cosmic rays heat and ionize the gas. Increasing $\chi_{\text{CR}}$ increases the densities of e$^-$, H$^0$, and H$^+$ in gas that is predominantly H$_2$. Further, H$^0$ and H$^+$ undergo exchange collisions with H$_2$ and induce ortho-para conversions. The upper half of Figure 3 shows how the column densities of various rotational levels of H$_2$ change with varying $\chi_{\text{CR}}$. This clearly illustrates that populations in the...
higher rotational levels of \( \text{H}_2 \) depend on \( \chi_{\text{CR}} \). The lower half of Figure 3 displays the column densities of\( \text{CO}, \text{OH}, \text{H}_2^+, \) and \( \text{NH} \). \( \chi_{\text{CR}} \geq 40 \) reproduces the observed column densities of the rotational levels of \( \text{H}_2^+(j = 0, 1, 2, 3, 4, 5) \), \( \text{H}_2^0, \text{CO}, \text{H}_3^+ \), and \( \text{NH} \), within a factor of 2 or better and reproduces the \( \text{H}_2 \) ortho-para temperature. However, the predicted column density of \( \text{OH} \) is too large for \( \chi_{\text{CR}} \geq 10 \). Interestingly, the \( \text{CO} \) and \( \text{OH} \) column densities change together as \( \chi_{\text{CR}} \) changes, suggesting that the problem is actually in the \( \text{CO}/\text{OH} \) balance. The dashed vertical line highlights \( \chi_{\text{CR}} = 40 \), the best-fit value.

3.3. Characteristics of the Best-fit Model

3.3.1. Overview of the Results

Black lines in Figure 1 plot the hydrogen ionization structure with \( n(\text{H}^0) \sim 35 \text{ cm}^{-3} \) and \( n(\text{H}^+) \sim 45 \text{ cm}^{-3} \) at the illuminated face for the \( \chi_{\text{CR}} = 40 \) model. These number densities at the illuminated face resemble the earlier model with \( \chi_{\text{CR}} = 1 \), since the shallow regions are photon dominated. Deep inside the cloud hydrogen is mainly in \( \text{H}_2 \) and \( f(\text{H}_2) \sim 0.78 \). The value of \( f(\text{H}_2) \) is smaller than the earlier model with \( \chi_{\text{CR}} = 1 \) due to enhanced cosmic-ray destruction of \( \text{H}_2 \). The temperature also increases deep inside the cloud compared to the \( \chi_{\text{CR}} = 1 \) model. Similarly, black lines in Figure 2 show that \( n_e \) and \( T_e \) vary from 45 to 0.034 cm\(^{-3} \) and 5370 to 51 K as functions of cloud depth, and clearly shows that the enhanced \( \chi_{\text{CR}} \) increases \( n_e \) and \( T_e \) deep inside the cloud. The fifth column of Table 3 lists the predicted column densities with \( \chi_{\text{CR}} = 40 \).

The merit of this model is that the enhanced \( \chi_{\text{CR}} \) not only reproduces the column densities of \( \text{H}_2^+ \), \( \text{H}_2^0 \), \( \text{CO} \), and \( \text{NH} \) within a factor of 2 but also the column densities of higher rotational levels of \( \text{H}_2 \) and the measured \( \text{H}_2 \) ortho-para temperature. In our calculation we use the dissociative recombination rate of \( \text{H}_2^+ \) measured by MC03. For a comparison, we also show the resulting \( \text{H}_2^+ \) abundance with a dissociative recombination rate of \( \text{H}_2^+ \) as derived by Datz et al. (1995) and tabulated by Stancil et al. (1998). The Datz et al. (1995) rate, being larger than that reported by MC03, produces a lower \( \text{H}_2^+ \) column density. Note that we do not require shocks to reproduce the column densities of the higher rotational levels of \( \text{H}_2 \). It is achieved by the cosmic-ray excitation of the ground state levels. Higher \( \chi_{\text{CR}} \) excites the electronic states more efficiently and enhances the populations of the higher vib-rotational levels of the ground state. Earlier van Dishoeck & Black (1986) also reproduced the column densities of the higher rotational levels of \( \text{H}_2 \) toward \( \zeta \) Persei without shocks by using an assumed radial gradient in temperature and density.

Our model reproduces the column densities of \( \text{H}_2^0, \text{O}_2^+, \text{CO}^+, \text{S}^0, \text{S}^+, \text{Si}^+, \text{H}_3^+, \) and \( \text{CO}, \text{NH} \) within a factor of 2. We reproduce the observed ratio of \( \text{S}^+/\text{S}^0, \text{Fe}^+/\text{Fe}^+, \text{C}^+/\text{C}^0, \text{Ca}^+/\text{Ca}^0 \), and the corresponding column densities within 2 \( \sigma \) of the observational uncertainties, which suggests that our predicted \( \chi/n_e \) is consistent with the observed data. Furthermore, the column densities determined from [C i] fine-structure excitations are reproduced verifying our predicted gas density.

3.3.2. \( \text{CH}^+ \) and \( \text{CH}^+ \)

Our calculation, like all those that use a standard ISM chemistry network, underpredicts the abundances of \( \text{CH} \) and \( \text{CH}^+ \). The basic problem is that the reactions forming these molecules have substantial energy barriers and do not occur at low temperatures. Investigations have suggested that these species can be produced by either shocks or nonthermal chemistry. Zsargó & Federman (2003) favor nonthermal chemistry. In the nonthermal chemistry driven by MHD (magnetohydrodynamical) waves, an effective temperature is defined based on the turbulence, which then allows the above reactions to occur. Shocks are another alternative. In this case the gas temperature is high and the reactions can occur.

If \( \text{CH}^+ \) is produced by shocks, \( \text{OH} \) will be produced as well. At high temperatures, \( \text{CH}^+ \) is formed by \( \text{C}^+ + \text{H}_2 \rightarrow \text{CH}^+ + \text{H} \) (reaction I) and \( \text{OH} \) by \( \text{O} + \text{H}_2 \rightarrow \text{OH} + \text{H} \) (reaction II). There is evidence suggesting that the formation of \( \text{CH}^+ \) is not always associated with the production of \( \text{OH} \) (for example, toward \( \zeta \) Per). Federman et al. (1996) showed that reaction I can occur without II if MHD waves only interact with the ions and this chemistry scheme can reproduce observed \( \mathcal{N}(\text{CH}^+) \) without producing excessive \( \mathcal{N}(\text{OH}) \). We have added the nonthermal chemistry caused by MHD waves as an option in Cloudy. We also included the ability to couple the MHD waves to all reactions, or to just the reactions where one of the reactants is an ion. Our predicted column densities of \( \text{CO}, \text{OH}, \text{CH}, \text{CH}^+, \text{CN}, \text{C}_2, \) and \( \text{C}_3 \) with nonthermal chemistry are listed in Table 3 (in the last column, marked by footnote c). We find reasonable agreement with the observed \( \mathcal{N}(\text{CH}^+) \) with little increase in \( \mathcal{N}(\text{OH}) \) for nonthermal chemistry.
The CO/OH ratio is almost independent of $\chi_{CR}$ (CO/OH $\sim$ 1.07 for $\chi_{CR} > 20$; see Fig. 3). Our predicted CO/OH ratio is $\sim$10 times smaller than the observed value, suggesting that the problem is in the chemistry rather than $\chi_{CR}$. In the ISM, the formation of OH starts with the creation of O$^+$ followed by ion-molecule reactions. CO is formed mainly via C$^+$ + OH $\rightarrow$ CO + H$^+$. The rate coefficient adopted for this process by Cloudy is taken from Dubernet et al. (1992). CO can also form from HCO$^+$. If we accept these values of the CO formation rate coefficients, then the CO/OH ratio depends primarily on the photodissociation rates of CO and OH. Our CO and OH photodissociation rates are taken from van Dishoeck (1988) and the UMIST database, respectively. While there has been considerable effort in the past to obtain reliable CO and OH photodissociation rates, some uncertainty may still remain. Recent studies have found that experimental oscillator strengths for photodissociation of CO are larger than those included in UMIST (Federman et al. 2003; Eidelsberg et al. 2006). This creates a faster dissociation rate when the gas is optically thin to the CO electronic lines, and more self-shielding when the lines become optically

### TABLE 3

**Column Densities of Various Species**

<table>
<thead>
<tr>
<th>Species</th>
<th>Observed log $N$ (cm$^{-2}$)</th>
<th>Predicted log $N$ (cm$^{-2}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>$\chi_{CR} = 1$</td>
</tr>
<tr>
<td>H$^0$ (1)</td>
<td>20.76–20.85</td>
<td>19.90</td>
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<tr>
<td>C$^0$ (2)</td>
<td>15.46–15.56</td>
<td>15.22</td>
</tr>
<tr>
<td>C$^+$ (2)</td>
<td>15.44–15.52</td>
<td>15.13</td>
</tr>
<tr>
<td>C$^+$ (2)</td>
<td>14.04–14.14</td>
<td>13.74</td>
</tr>
<tr>
<td>CO$^+$ (4)</td>
<td>9.70–9.82</td>
<td>9.37</td>
</tr>
<tr>
<td>O$^+$ (5)</td>
<td>13.12–13.52</td>
<td>13.20</td>
</tr>
<tr>
<td>Ar$^+$ (2)</td>
<td>14.65–15.01</td>
<td>15.20</td>
</tr>
<tr>
<td>Mg$^+$ (2)</td>
<td>13.27–13.67</td>
<td>13.38</td>
</tr>
<tr>
<td>Fe$^+$ (4)</td>
<td>11.70–11.86</td>
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<tr>
<td>Fe$^+$ (4)</td>
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<td>Na$^+$ (4)</td>
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<td>C$_2$ (14)</td>
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<td>C$_3$ (15)</td>
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</tr>
<tr>
<td>H$_2$(J = 0) (2)</td>
<td>20.34–20.68</td>
<td>20.73</td>
</tr>
<tr>
<td>H$_2$(J = 1) (2)</td>
<td>20.01–20.36</td>
<td>20.33</td>
</tr>
<tr>
<td>H$_2$(J = 2) (2)</td>
<td>18.02–18.08</td>
<td>17.56</td>
</tr>
<tr>
<td>H$_2$(J = 3) (2)</td>
<td>16.11–16.43</td>
<td>15.55</td>
</tr>
<tr>
<td>H$_2$(J = 4) (2)</td>
<td>14.95–15.07</td>
<td>14.45</td>
</tr>
<tr>
<td>H$_2$(J = 5) (2)</td>
<td>14.32–14.40</td>
<td>14.06</td>
</tr>
<tr>
<td>T$_{10}$ (K)</td>
<td>45–79</td>
<td>~54</td>
</tr>
</tbody>
</table>

* Using recombination rate of H$_2^+$ measured by MC03.
* Using recombination rate of H$_2^+$ measured by Datz et al.
* Using nonthermal chemistry.
* Using 3/2 times decreased photodissociation rate.
* Using 3 times increased photodissociation rate.

thick. Experimental oscillator strengths are not available for the majority of the CO electronic lines, establishing an uncertainty in the predictions. Improved radiative rates for CO along with a detailed treatment of the radiative transfer using more modern grain physics are required. In the case of OH, the photodissociation cross sections are highly accurate, but standard, and potentially out-of-date ISM grain models were used to deduce attenuated photodestruction rates. Further, these rates were fitted to simple functions which may culminate in considerable uncertainty, up to nearly a dex. For instance, a factor of 3 increase in the photodissociation rate of OH and a factor of 3/2 decrease in the photodissociation rate of CO results in column densities that are within a factor of 2 of the observed range. If these uncertainties are considered the differences are not serious.

As mentioned above, we do not try to reproduce the molecules that form in a dense molecular region (C$_2$, C$_3$, and CN) or the highly ionized species that form in the H$^+$ region near the background star (S$^{+2}$ and Si$^{+2}$). We do list our predictions of these species for the diffuse phase.

For comparison we list our predicted column densities for $\chi_{\text{CR}} = 10$ in the fourth column of Table 3. Our predicted OH column density matches the observation better for $\chi_{\text{CR}} = 10$ but the predicted column densities of the other species (H$^0$, O$^0$, C$^0$, H$_2^+$, CO, NH, and higher rotational levels of H$_2$) are lower than the observed values. In contrast, we reproduce the observed column densities within a factor of 2 for $\chi_{\text{CR}} = 40$ with the one exception of OH. We cannot decrease $\chi_{\text{CR}}$ in order to reproduce OH because in that case the model will underpredict other species, including the high-$J$ excitation in H$_2$, by more than a factor of 3. Keeping in mind the possible uncertainties in the photodestruction rates of OH and CO, the $\chi_{\text{CR}} \sim 40$ case is certainly a possible solution in the diffuse cloud toward $\zeta$ Persei.

### 3.3.4. The Grain-size Distribution

Figure 4 and Table 4 show how details of the grain physics affect the H$_2^+$ column density and so the deduced $\chi_{\text{CR}}$. We perform tests with different grain-size distributions, keeping all other parameters fixed. The figure shows the ratio $n$(H$_2^+)/n$(H$_2$), which is proportional to $\chi_{\text{CR}}/n_e$ (see eq. [1]). This sight line is characterized by $R_s = 2.85$, and we use a custom size distribution with a greater than standard number of small particles to approximate the extinction on this line of sight. This will increase the rate of charge exchange with ions, with small grains capturing free electrons, and absorbing FUV radiation. Since $\chi_{\text{CR}}$ is held fixed, the changes shown in Figure 4 are mainly due to changes in $n_e$. Had we assumed, like many ISM studies (for example MC03), that the electron density is equal to the carbon density, we would have needed a value of $\chi_{\text{CR}}$ almost 2 times larger to reproduce $N$(H$_2^+)/N$(H$_2$). To see the effect of collisional grain physics for this sight line we turned off grain-ion charge exchange as well as electron capture by grains. As a result $n_e$ increases, mainly because elements with low ionization thresholds stay ionized and donate more electrons to the gas. This results in a lower value of $N$(H$_2^+)/N$(H$_2$). We find that grain-ion charge exchange plays an important role in this model.

The temperature is derived from the balance between heating and cooling processes. The gas heating occurs mainly via grain photoelectric heating, cosmic-ray heating, heating due to H$_2$ dissociation, and collisional de-excitation, whereas the gas cooling occurs mainly via collisionally excited fine-structure atomic and molecular lines. In Table 3 we list the excitation temperature ($T_{\text{ex}}$) derived from the $J = 1$ and $J = 0$ levels of H$_2$. Deep inside the cloud, cosmic rays also heat the gas, and an increase in $\chi_{\text{CR}}$ will increase the temperature (and hence increase $T_{\text{ex}}$). For $\chi_{\text{CR}} = 40$, the cosmic-ray heating rate deep inside the cloud is $\sim 9.3 \times 10^{-25}$ erg cm$^{-3}$ s$^{-1}$, which is roughly half of the total heating.

S06 noticed that $\chi_{\text{CR}} = 20$ increased $T_{\text{ex}}$ by a factor of 2 above $\chi_{\text{CR}} = 1$ along the line of sight toward HD 185418. However, toward $\zeta$ Persei the temperature did not increase significantly even though we have a higher $\chi_{\text{CR}} = 40$. One contributor is the smaller $R_s = 2.85$ and hence the presence of more small grains toward $\zeta$ Persei. For a given mass, the smaller grains have more
surface area and produce more heating compared to larger grains. As a result, deep inside the cloud the heating due to cosmic rays is a smaller fraction of the total. Furthermore, the density along the \( \zeta \) Persei line of sight is higher. Since the effects of cosmic-ray heating are proportional to \( \chi_{\text{CR}}/n_\text{H} \), the larger \( \chi_{\text{CR}} \) is compensated by the larger \( n_\text{H} \).

### 3.3.5. Unobserved Species

We list the predicted column densities for some unobserved species in Table 5. The predicted column densities of \( \text{OH}^+ \) and \( \text{H}_2^+ \) for \( \chi_{\text{CR}} = 40 \) are \(~4\) times larger than that for \( \chi_{\text{CR}} = 10 \). Future detection of these species will help to distinguish between these two cases. It should be possible to measure hydrogen recombination lines along this sight line, but these would form in the \( \text{H}^+ \) region that produces \( \text{S}^- \) and \( \text{Si}^- \) rather than in cosmic-ray ionized gas. In our model high-\( J \) levels of \( \text{H}_2 \) are excited by cosmic rays. If instead these excitations are produced in shocks then one will be able to detect the differences in the velocity of the absorption line using high-resolution spectroscopic data. Absence of any such velocity difference will support the models presented here.

Recently Padoan & Scalo (2005) have shown that self-generated MHD waves produce an enhanced cosmic-ray density in diffuse clouds compared to those found in dense clouds. They have shown the variation of the cosmic-ray density as a function of \( n_\text{H} \) in their Figure 1 and suggest a scaling law such as \( \chi_{\text{CR}} \propto n_\text{H}^{0.5} \), where \( n_\text{H} \) is the ion density. The density of the diffuse cloud toward \( \zeta \) Persei lies near the top of the curve predicting \( \chi_{\text{CR}} \). In an earlier paper (S06) we required 20 times \( \chi_{\text{CR}} \) to reproduce the observed column density along a sight line for which \( n_\text{H} = 27 \) cm\(^{-3} \). Assuming \( \chi_{\text{CR}} \propto n_\text{H}^{0.5} \) and rescaling the \( \chi_{\text{CR}} \) found by S06 by the ratio of densities according to the Padoan & Scalo prescription we should get \(~34.5 \) times \( \chi_{\text{CR}} \) in the diffuse cloud toward \( \zeta \) Persei, which lies very close to the lower limit of our range of suggested \( \chi_{\text{CR}} \).

### Table 4

<table>
<thead>
<tr>
<th>Species</th>
<th>Column Densities with Larger Grains (Orion Type, ( R_I = 5.5 ))</th>
<th>Column Densities with Smaller Grains (ISM Type, ( R_I = 3.1 ))</th>
<th>Observed Column Densities for ( \zeta ) Per</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \text{H}_2^+ )</td>
<td>12.83</td>
<td>13.50</td>
<td>14.08</td>
</tr>
<tr>
<td>( \text{H}_2 )</td>
<td>20.36</td>
<td>20.66</td>
<td>20.78</td>
</tr>
<tr>
<td>( \text{H}_2^+ )</td>
<td>12.88</td>
<td>13.18</td>
<td>13.34</td>
</tr>
<tr>
<td>( \text{S}^0 )</td>
<td>13.55</td>
<td>13.65</td>
<td>13.59</td>
</tr>
<tr>
<td>( \text{S}^+ )</td>
<td>16.30</td>
<td>16.30</td>
<td>16.30</td>
</tr>
<tr>
<td>( \text{C}^+ )</td>
<td>15.57</td>
<td>15.70</td>
<td>15.60</td>
</tr>
<tr>
<td>( \text{C}^0 )</td>
<td>17.30</td>
<td>17.30</td>
<td>17.30</td>
</tr>
<tr>
<td>( \text{Fe}^+ )</td>
<td>11.69</td>
<td>11.63</td>
<td>11.66</td>
</tr>
<tr>
<td>( \text{Fe}^2+ )</td>
<td>14.20</td>
<td>14.20</td>
<td>14.20</td>
</tr>
</tbody>
</table>

### Table 5

<table>
<thead>
<tr>
<th>Species</th>
<th>Predicted log ( N ) cm(^{-2} ) ( (\chi_{\text{CR}} = 10) )</th>
<th>Predicted log ( N ) cm(^{-2} ) ( (\chi_{\text{CR}} = 40) )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \text{H}^+ )</td>
<td>18.10</td>
<td>18.19</td>
</tr>
<tr>
<td>( \text{H}_2^+ )</td>
<td>12.84</td>
<td>13.38</td>
</tr>
<tr>
<td>( \text{OH}^+ )</td>
<td>12.16</td>
<td>12.79</td>
</tr>
<tr>
<td>( \text{H}_2^+ )</td>
<td>12.03</td>
<td>12.54</td>
</tr>
<tr>
<td>( \text{H}_2O^+ )</td>
<td>12.01</td>
<td>12.22</td>
</tr>
</tbody>
</table>

### 4. CONCLUSIONS

The observation of a large \( \text{H}_2^+ \) column density toward \( \zeta \) Persei raises a controversy about the value of cosmic-ray ionization rate toward this line of sight. Earlier McCall et al. (2003) derived an enhanced \( \chi_{\text{CR}} = 48 \) toward \( \zeta \) Persei by assuming \( n(\text{H}) = 250 \) cm\(^{-3} \) and that most of the electron density \( n_e \) comes from carbon photoionization. On the other hand, LP04 found that \( \chi_{\text{CR}} = 10 \) can reproduce almost all observed column densities within a factor of 3 based on a detailed model consisting of two thermal phases plus three shocks. However, our detailed model reproduces many of the observed column densities in the diffuse cloud toward \( \zeta \) Persei with the help of an enhanced cosmic-ray ionization rate. We find that with the exception of \( \text{OH} \), an enhancement of \( \chi_{\text{CR}} \geq 40 \) can reproduce the observed column densities of \( \text{H}_2^+ \), \( \text{H}_3^0 \), \( \text{CO} \), \( \text{NH} \), and various \( \text{H}_2 \) lines, within a factor of 2. The merit of our model is that we do not require shocks to reproduce the observed \( N(\text{CH}^+) \) and populations of high rotational levels of \( \text{H}_2(J = 2, 3, 4, \text{ and } 5) \). We use nonthermal chemistry driven by MHD waves, which has an advantage of producing the observed \( N(\text{CH}^+) \) without producing excessive \( \text{OH} \). The high rotational levels of \( \text{H}_2 \) are produced as a result of excitation by cosmic rays. The value of the cosmic-ray ionization rate is model dependent since some assumptions come into play. Tests presented above show that the grain physics has a significant effect on the deduced \( \chi_{\text{CR}} \). Small grains, which are overabundant along this low-\( R_I \) sight line, neutralize ions, remove free electrons, and extinguish FUV radiation. These effects by themselves change the deduced \( \chi_{\text{CR}} \) by more than a factor of 2. This raises the question of the composition of the small grains; this is uncertain, with some favoring PAHs rather than graphite, as we assume. The main unknown, however, is the grain size distribution itself, which introduces a further factor of 2 uncertainty in addition to the unknowns presented by the chemical network and cloud geometry.

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