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C. R. O'Dell
Vanderbilt University

W. Kollatschny
Universität Göttingen, Germany

Gary J. Ferland
University of Kentucky, gary@uky.edu

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Which Stars Are Ionizing the Orion Nebula?

C. R. O’Dell¹, W. Kollatschny², and G. J. Ferland³

¹ Department of Physics and Astronomy, Vanderbilt University, Nashville, TN 37235-1807, USA

² Institut für Astrophysik, Universität Göttingen, Friedrich-Hund-Platz 1, D-37077 Göttingen, Germany

³ Department of Physics and Astronomy, University of Kentucky, Lexington, KY 40506, USA

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Abstract

The common assumption that θ^1 Ori C is the dominant ionizing source for the Orion Nebula is critically examined. This assumption underlies much of the existing analysis of the nebula. In this paper we establish through comparison of the relative strengths of emission lines with expectations from Cloudy models and through the direction of the bright edges of proplyds that θ^2 Ori A, which lies beyond the Bright Bar, also plays an important role. θ^1 Ori C does dominate ionization in the inner part of the Orion Nebula, but outside of the Bright Bar as far as the southeast boundary of the Extended Orion Nebula, θ^2 Ori A is the dominant source. In addition to identifying the ionizing star in sample regions, we were able to locate those portions of the nebula in 3D. This analysis illustrates the power of MUSE spectral imaging observations to identify sources of ionization in extended regions.

Key words: H II regions – ISM: individual objects (Orion Nebula) – stars: individual (θ^1 Ori C, θ^2 Ori A)

1. Introduction

The central region of the Orion Nebula (NGC1976) is arguably the best-studied H II region (O’Dell et al. 2008). As the exemplar of its class of gaseous nebulae, it has also been subject to many attempts at modeling the physics that occurs and the 3D structure. We understand that it is basically an irregular concave thin layer of ionized gas lying nearer the observer than the main ionization front (MIF) that marks the boundary with the host background molecular cloud. There is also a nearby foreground thin layer (the Veil) of partially ionized gas.

The usual assumption is that ionization in this region is dominated by the brightest star in the compact Trapezium group, the complex hot star θ^1 Ori C, whose spectral type is usually assigned as about O6. The next hottest and most luminous star is θ^2 Ori A, which lies $135''$ southeast of θ^1 Ori C. As one moves farther from the Trapezium, the degree of ionization decreases, consistent with the ionizing star or stars being in the central nebula (O’Dell & Harris 2010). It is important in interpreting the nebula’s spectra to understand whether a secondary star or stars play a role in ionization of the nebula.

The inner ionized nebula is actually very complex in structure, there being irregular features in the concave surface, the most famous of which is the Bright Bar, and there is a neutral cloud of material containing very young stars known as the Orion-South cloud (hereafter Orion-S) lying $55''$ to the southwest of θ^1 Ori C. Interpreting the spectra and emission-line images of the nebula demands knowing the source of ionization for each region. That is the goal of this study.

Two approaches are adopted. In Section 2 we utilize the predictions of key tracers of different degrees of ionization. In Section 3 we utilize resolved ionization fronts formed around gas surrounding many of Orion’s proplyds.

Numerous assumptions have been made in this study. There are many similar values for the distance to the Orion Nebula, but in this study we have adopted 388 ± 8 pc from the recent radio results of Kounkel et al. (2017). For θ^1 Ori C, we have adopted a temperature (T_{star}) of 38,950 K and total luminosity in photons capable of ionizing hydrogen (Q_{LyC}) of 7.35×10^{48}

photons s^{-1} (Badnell et al. 2015). For θ^2 Ori A at O9.5 V (Warren & Hesser 1977), we use values from Table 2.3 of Osterbrock & Ferland (2006) of $T_{\text{star}} = 34,600$ K and $Q_{\text{LyC}} = 3.63 \times 10^{48}$ photons s^{-1} . We express distances in the plane of the sky in parsecs, using the adopted distance.

2. Using Line Ratios to Identify the Ionizing Star

The ionized layer on the observer’s side of the MIF is stratified into layers according to the varying ionization. These layers are determined by the energy of the stellar photons reaching them, and this energy distribution is controlled by the absorption of the ions of the most common elements, hydrogen and helium.

Close to the ionizing star hydrogen is ionized and oxygen is doubly ionized (easily traced by the [O III] 500.7 nm line). Helium is singly ionized (easily traced by the He I 667.8 nm line). This is the $\text{He}^+ + \text{H}^+$ zone. For the relatively cool Orion Nebula stars there is no higher ionization zone (where helium would be doubly ionized).

Farther from the ionizing star and closer to the MIF is the narrow $\text{He}^0 + \text{H}^+$ zone, where helium is neutral and hydrogen is ionized. Nitrogen is singly ionized and is most visible in the [N II] 658.3 nm line. The easily visible Balmer $\text{H}\alpha$ and $\text{H}\beta$ lines arise from both ionization zones.

2.1. The Basic Approach

We selected seven regions that were expected to illuminate the question of the ionizing star in various parts of the Orion Nebula. These all fall within the region of calibrated monochromatic images obtained with the MUSE (Weilbacher et al. 2015) multi-aperture spectrograph. Reddening-corrected line ratios from the MUSE database are used throughout this study. Figure 1 shows the location and sizes of these samples. They include regions near θ^1 Ori C and others on both sides of the Bright Bar.

The line ratios $\text{F}([\text{O III}])/\text{F}(\text{H}\beta)$ and $\text{F}(\text{He I})/\text{F}(\text{H}\alpha)$ within the $\text{He}^+ + \text{H}^+$ zone and $\text{F}([\text{N II}])/\text{F}(\text{H}\alpha)$ within the $\text{He}^0 + \text{H}^+$ zone are dependent on the temperature of the ionizing star and a quantity defined as the ionization parameter U (Osterbrock &

Table 1
Ionization Parameters^a and Flux Ratios of the Seven Samples

Sample	n_e ^b	$\log U(\theta^1 \text{ Ori C})$	$\log U(\theta^2 \text{ Ori A})$	$F([\text{O III}] 500.7 \text{ nm})/F(\text{H}\beta)$	$F([\text{N II}] 658.3 \text{ nm})/F(\text{H}\alpha)$	$F(\text{He I } 667.8 \text{ nm})/F(\text{H}\alpha)$
A	6160	-0.95	-2.92	3.73	0.120	0.0129
B	4310	-0.24	-2.54	3.58	0.136	0.0126
C	2090	-1.56	-1.28	2.81	0.125	0.0132
D	1230	-1.76	-0.81	1.06	0.307	0.00812
E	1610	-1.71	-1.26	1.96	0.180	0.0118
F	1860	-1.58	-1.22	2.93	0.116	0.013
G	1970	-1.27	-1.70	3.03	0.129	0.0127

Notes.

^a The ionization parameters are upper limits, being calculated from the distance to the ionizing star as projected on the sky.

^b Electron densities are in cm^{-3} and are derived from the red [S II] doublet ratio using MUSE fluxes.

Ferland 2006), which is the ratio of the density of photons that can ionize hydrogen to the local hydrogen density (n_{H}). We can approximate the local hydrogen density as the electron density (n_e) since hydrogen is essentially completely ionized and the $\sim 10\%$ contribution from He I is small.

We can calculate the photon density (ρ_{LyC}) from the total ionizing photon luminosity of the ionizing star (Q_{LyC}) and the true distance to the emitting layer (r), and this leads to the relation

$$U = \frac{Q_{\text{LyC}}}{4\pi r^2 c} \times \frac{1}{n_e}, \quad (1)$$

where c is the velocity of light.

2.2. Calculations

In this section we present a series of new ionization simulations of the Orion H II region. We use version C17.00 of Cloudy, the spectral simulation code last described by Ferland et al. (2013). The geometry is that used by Baldwin et al. (1991). It is assumed that the physical thickness of the H^+ layer that constitutes the H II region is much thinner than the separation to the ionizing stars. The geometry is plane parallel. The layer is assumed to be in hydrostatic equilibrium with the H^+ layer held back against the molecular cloud by the combination of gas pressure in the X-ray-emitting hot gas surrounding the star cluster and the absorbed outward momentum in the starlight.

This model was originally applied to central regions of the nebula, close to the Trapezium, where we are most likely viewing the H^+ layer nearly face-on. However, the bright Huygens region is most likely bowl shaped, as described by Ferland et al. (2013) and Wen & O'Dell (1995). The observed surface brightness of the layer is affected by two things. First, as the distance of the star increases, the flux of ionizing photons, which sets the surface brightness, falls as r^{-2} . This causes the surface brightness to decrease with increasing radius. At the same time, the “tilt” of the H^+ layer to our line of sight increases (see Figure 3 of Wen & O'Dell 1995), which causes the surface brightness to *increase* by a factor of $\cos(\theta)^{-1}$, where θ is the viewing angle measured relative to the normal to the layer. As long as our line of sight passes through the entire layer, the relative emission-line intensities do not change, only their total surface brightness.

This simple geometry breaks down when the viewing angle approaches edge-on, since we then see successive regions projected on the plane of the sky, resolving the ionization structure. This occurs within the Bright Bar (Sellgren et al. 1990)

or near edges of proplyds. In this case a more complex model can be developed (Shaw et al. 2009). For simplicity, in this study we use the Baldwin et al. (1991) approach and remain mindful that it will break down in edge-on cases.

We calculated the line ratios $F([\text{O III}], 500.7 \text{ nm})/F(\text{H}\beta)$, $F(\text{He I } 667.8 \text{ nm})/F(\text{H}\alpha)$, and $F([\text{N II}], 658.3 \text{ nm})/F(\text{H}\alpha)$ for a variety of conditions. These assumptions are summarized in Table 1. The models adopted a progression of values of the ionizing star temperature and the ionization parameter. We used the stellar spectral energy distribution described in Badnell et al. (2015). In Figures 2 and 3 the predicted values of the line ratios are plotted as contour lines. The lines in these ratios are at close wavelengths and relatively insensitive to reddening by foreground dust. Furthermore, the models were compared with reddening-corrected line ratios.

2.3. Methodology

By comparing the derived data for each sample with the expected values for the candidate ionizing star, we can identify that star. The apparent value of U (U_{sky}) was calculated from the equation

$$U_{\text{sky}} = \frac{Q_{\text{LyC}}}{4\pi r_{\text{sky}}^2 c} \times \frac{1}{n_e}, \quad (2)$$

where r_{sky} is the separation of the sample and the candidate ionizing star and we have adopted the n_e value derived from the red [S II] ratio. This calculation was done for both $\theta^1 \text{ Ori C}$ and $\theta^2 \text{ Ori A}$. This means that a given sample is represented twice, once in Figure 2 and again in Figure 3, the first using U_{sky} for $\theta^1 \text{ Ori C}$ and the second using U_{sky} for $\theta^2 \text{ Ori A}$, but with the same line ratio for the sample.

We do not know r , but we do know a lower limit, r_{sky} . Therefore, an ionization parameter calculated using r_{sky} (U_{sky}) is an upper limit. If the assumed ionizing star is correct, then each derived point will lie on the intersection of the observed line ratio and T_{star} or to the upper left. A greater separation means a larger difference between r and r_{sky} , that is, there must be a greater distance correction.

If the observed point lies to the right of the intersection or if a projection along a constant line ratio does not reach the T_{star} value, then that sample is not ionized by the assumed star. We do see two points (C and F) in the $F([\text{N II}])/F(\text{H}\alpha)$ panel of Figure 2 that lie to the right of the intersection with T_{star} for $\theta^1 \text{ Ori C}$. These fall in a very flat region for the predicted flux ratio. If we had adopted a smaller distance to the Orion Nebula, then the values of U_{sky} would be larger for all points, and Samples C and F would have

Figure 1. This $H\alpha$ 0.66×0.55 pc ($353'' \times 295''$) image of the Orion Nebula is from the $0''.2$ pixel $^{-1}$ MUSE data. The samples employed in generating the data in Table 1 are shown as lettered rectangles the size of the sample. The arrows indicate the direction(s) of the orientation of the bright edges of some of the many proplyds, as discussed in Section 3. Those proplyds ionized by θ^2 Ori A are shown in enlargements of $2'' \times 2''$, except for 244.3–439.8, which is a sample of $4'' \times 4''$. The location of the first of the *Spitzer Space Telescope* observations discussed in Section 4 is shown south of θ^2 Ori A and further defined in Figure 4.

shifted to near the intersection. However, our adopted value for the distance is already smaller than many modern determinations, as summarized in O'Dell & Henney (2008).

It is more likely that the Cloudy model that we adopt is beginning to break down in Samples C and F as a result of the large viewing angle. Our model assumes that we are observing along a line of sight that goes through the entire H^+ layer. When the MIF is highly tilted, we resolve the ionization structure on the plane of the sky and thus sample along layers within the column. This compromises a comparison of line-of-sight models against observations of a highly tilted region. We would expect that the results for the $F([N II])/F(H\alpha)$ ratio would be affected the most since the $[N II]$ line forms in a small part of the H^+ layer. Although the peak of the tilt of the MIF is a maximum at the nearby Bright Bar, O'Dell et al. (2017) establish that the region near our Samples C and F is already highly tilted.

2.4. Identifying the Ionizing Star from the Observed Ratios

In Figure 2 we compare our observations with the brightest star in the Orion Nebula, θ^1 Ori C. Figure 3 is similar, but assumes that the isolated bright star θ^2 Ori A is the source.

In Figure 2 we see that Samples A, B, C, F, and G are all ionized by θ^1 Ori C, although the $F([N II])/F(H\alpha)$ ratios for Samples C and F fall to the right of $T_{\text{star}}(\theta^1 \text{ Ori C})$. In

$F([O III])/F(H\beta)$ and $F(He I)/F(H\alpha)$ we see that all samples are consistent with ionization by θ^1 Ori C, with the caveat of the remarks in Section 2.3 about the high T_{star} values for Samples C and F. Taken at face value, $F([O III])/F(H\beta)$ and $F(He I)/F(H\alpha)$ indicate that Samples C and F have spatial distances about equal to r_{sky} , while Samples G, A, and B demand progressively larger corrections to the true distance. In the left and right panels it appears that Sample D could also be associated with θ^1 Ori C, but in the middle panel we see that no plausible correction for distance would make the line ratio compatible with T_{star} for θ^1 Ori C. All of the θ^1 Ori C associated samples fall within the central cavity of the nebula. Sample E can be associated with ionization by θ^1 Ori C with a small adjustment in U, which is allowable even though it is much further on the sky than the difficult-to-interpret Samples C or F.

In Figure 3 we note that only Sample D agrees with ionization by θ^2 Ori A with similar distance corrections in all ratios. The appearance of the ionized edge of the nearby proplyd 244.3–439.8 indicates dominant ionization by θ^2 Ori A, but with some possible contribution by θ^1 Ori C (Section 3). Clearly Samples A, B, C, F, and G are not associated with θ^2 Ori A. Those samples lie on the θ^1 Ori C side of the Bright Bar and have already been associated with that star. Sample E agrees with an association with θ^2 Ori A in the $F([O III])/(H\beta)$ ratio without a distance correction, but in the $F(He I)/F(H\alpha)$ and $F([N II])/F(H\alpha)$

Figure 2. Ionization parameter vs. ionizing star temperature is shown in all three panels, each depicting extinction-insensitive line ratios. Dashed vertical lines indicate the temperatures of the two candidate ionizing stars. The dotted lines indicate the expected line ratio using Cloudy for varying values of the distance-dependent ionization parameter assuming that the ionizing star is θ^1 Ori C. The seven samples from Figure 1 are shown. Those probably ionized by θ^1 Ori C are in blue and circled. Arrows indicate the direction of the correction to the ionization parameter for the fact that it was calculated using the minimum possible distance from the star, i.e., the distance in the plane of the sky. The length of the arrow only suggests the magnitude of that correction. Blue circles and lettering indicate samples most likely ionized by θ^1 Ori C.

ratios it lies too far to the right of T_{star} for θ^2 Ori A, which rules out θ^2 Ori A.

The identification of the inner five samples as being ionized by θ^1 Ori C falls within the accepted model of this region of the Orion Nebula, i.e., that it is a concave irregular surface. The sample apparently ionized by θ^2 Ori A is consistent with the Bright Bar and the regions southeast from there being ionized by the nearest hot star. Sample E, although lying outside the Bright Bar, is probably ionized by θ^1 Ori C, as discussed earlier in this section.

3. Identifying the Ionizing Stars Using Observations of the Proplyds

There is a completely independent method of identifying the dominant source of ionization using high-resolution images of the Orion Nebula proplyds. These are young stars with circumstellar material in or projected on H II regions. Under favorable orientations and locations the interior protoplanetary disks are rendered visible in silhouette. More relevant here, these disks are surrounded by much larger envelopes of gas and dust. A nearby hot star will form an ionization front in this gas that is seen as a bright arc oriented toward the ionizing star. This allows one to identify the direction in the plane of the sky of the dominant ionizing star.

In Figure 1 we have added white arrows on a representative sample of the proplyds. The arrows point in the direction of the bright arc pointing toward that proplyd's ionizing star. The sampling is incomplete, but representative of the proplyds closest to θ^1 Ori C, and includes all of the proplyds that indicate ionization by θ^2 Ori A. The latter four objects (O'Dell et al. 1993; O'Dell & Wen 1994) are labeled with their names in the position-based designation system introduced by O'Dell & Wen (1994) and later refined (O'Dell et al. 2015). The orientations of the bright arcs of those proplyds ionized by θ^1 Ori C are well illustrated in O'Dell (1998). Those four associated with θ^2 Ori A are shown in enlargements. It is to be noted that the circumstellar gas of 196.6–426.6 is clearly ionized by both stars, as may be the case for 244.3–439.8.

4. Using *Spitzer Space Telescope* Infrared Spectra Line Ratios

Published *Spitzer Space Telescope* observations of the infrared [Ne III] 15.6 μm and [Ne II] 12.8 μm lines are also useful for determining the ionizing star. These are particularly useful because they both arise in the same $\text{He}^+ + \text{H}^+$ zone and are unaffected by emission from the $\text{He}^0 + \text{H}^+$ zone and any changes of abundances, as is the case for our optical line ratios. In the study of Rubin et al. (2011) observations were made at 11 locations to the southeast of the central Orion Nebula as shown in Figure 4. These samples all lie outside the Bright Bar, with the innermost being close to θ^2 Ori A. The progression of increasing distances is I4, I3, I2, I1, M1, M2, M3, M4, V1, V2, and V3. They all lie in a smooth region between the Bright Bar and the boundary of the Extended Orion Nebula (EON). Sample V2 falls near the edge of the sharp southeast boundary of the EON, and Sample V3 falls on this boundary. These samples do not agree with the results from the others because Sample V3 is certainly an ionization front viewed edge-on and Sample V2 is probably a region that is rising toward that feature (Rubin et al. 2011), and thus our Cloudy models do not provide a guide to the conditions (Section 2.3). They are not similar to the other samples and are not included in our discussion of the general properties of the *Spitzer* samples.

We have compared the neon line ratios with the predictions from the same set of Cloudy calculations used in the analysis of the optical line ratios. We have used the Rubin et al. (2011) [S II] n_e densities derived from overlapping optical observations for each sample. The results are shown in Figure 5.

Figure 5 shows that all of the samples that assume θ^1 Ori C to be the ionizing star (blue labels) cluster near $T_{\text{star}} \sim 34,000$ K. This is clear evidence that θ^1 Ori C does not ionize the region southwest of the Bright Bar. To reconcile the observations with T_{star} for θ^1 Ori C would require all of the samples to have a large and similar correction for projection effects, even though their separations in the plane of the sky range from 2'6 to 12'08.

The comparison with the results assuming θ^2 Ori A (red labels) is fully consistent with that star as the dominant ionizing

Figure 3. Same as Figure 2, except now the ionization parameters are calculated using distance from θ^2 Ori A. The samples probably ionized by θ^2 Ori A are in red and circled.

Figure 4. *Spitzer* samples (after Figure 1 of Rubin et al. 2011); see text for details. The east boundary of the EON appears at the lower left (southeast), and the Trapezium stars lie in the center of the bright northwest feature.

star, since Sample I4 is closest to θ^2 Ori A and our Sample D in the plane of the sky and the I3, I2, and I1 samples progress farther away. Their locations in Figure 5 indicate a succession of corrections for projection effects, and at large distances the samples are close to T_{star} for θ^2 Ori A.

5. Discussion

The results of Sections 2 and 3 show similar but not identical results. Some of the differences may be explained by the methodology.

The $F([\text{O III}])/F(\text{H}\beta)$ and $F(\text{He I})/F(\text{H}\alpha)$ line ratios are indicators of the conditions in the $\text{He}^+ + \text{H}^+$ zone, as $[\text{O III}]$ and He I emission arises only there, as does most of the $\text{H}\alpha$ emission. The gas density decreases (O'Dell 2001) with increasing distance from the MIF; therefore, most of the emission used in the $F([\text{O III}])/F(\text{H}\beta)$ and $F(\text{He I})/F(\text{H}\alpha)$ line ratio method occurs near the onset of the $\text{He}^+ + \text{H}^+$ zone, with

Figure 5. Like Figures 2 and 3, except now we use the ratio of the $[\text{Ne III}]$ 15.6 μm and $[\text{Ne II}]$ 12.8 μm lines. The symbols are samples from the study of Rubin et al. (2011) that begin close to θ^2 Ori A and proceed to the southeast with a maximum distance of 12.08' from θ^1 Ori C. Colors indicate which star was assumed to be the ionizing source, blue for θ^1 Ori C and red for θ^2 Ori A. Samples V2 and V3 fall at and near the edge of the sharp southeast boundary of the EON and are not used in the discussion in Section 4.

a lesser $\text{H}\alpha$ contribution from the narrower $\text{He}^0 + \text{H}^+$ zone. The distributions of doubly ionized oxygen (producing the $[\text{O III}]$ emission) and He I ions (producing the 667.8 nm emission) are slightly different within the $\text{He}^+ + \text{H}^+$ zone. This means that the ratios using these two ions will give slightly different results. In the $\text{He}^0 + \text{H}^+$ zone the $[\text{N II}]$ emission is well constrained, but the fact that the $\text{H}\alpha$ emission comes mostly from the $\text{He}^+ + \text{H}^+$ zone means that the $F([\text{N II}])/F(\text{H}\alpha)$ ratio is intrinsically not as good a diagnostic.

In contrast, the proplyds give an indication of the ionization of objects located somewhere along the line of sight from the foreground Veil and the MIF. The Orion Nebula Cluster is centered near the Trapezium and monotonically decreases in

stellar density with increasing distance in all directions. This means that a proplyd lying in the same direction as a ratios sample can lie well displaced toward the observer. This difference of sampling can explain the similarities and differences in the results.

5.1. Location of the Samples along the Line of Sight

We can, however, approximately determine the position of the MIF samples with respect to the plane in the sky of the ionizing star. For each observed sample we can compare the “observed” U_{sky} with the value where the same line ratio intersects T_{star} . The shift in U gives us the square of r/r_{sky} , hence that ratio. For example, if U_{sky} is four times larger than where that line ratio intersects T_{star} , then U_{sky} is four times too large and the true distance (r) is twice the separation in the plane of the sky (r_{sky}). It follows that the angle of r with respect to the plane of the sky (Θ) will be $\arccos(r_{\text{sky}}/r)$ and the distance along the line of sight (Z) is $Z = r_{\text{sky}} \times \tan \Theta$. What we actually determine is $\pm\Theta$ and $\pm Z$.

Assignment of positive (toward the observer) or negative can usually be determined by comparing the Z values with respect to the quantitative models of the Orion Nebula. From the ratio of the surface brightness in the radio continuum with that expected from ionization by θ^1 Ori C, both Wen & O'Dell (1995) and Henney et al. (2005) derived a model of a concave surface with ridges (e.g., the one producing the Bright Bar). The former presumed a distance to the nebula of 500 pc and the latter 430 pc. Their derived separation (Z) values were -0.2 pc and about -0.15 pc. In their study of the spectrum of the substellar point using Cloudy, Badnell et al. (2015) used 0.10 pc. We assume in our discussion that the sub- θ^1 Ori C distance is -0.15 pc. The line-of-sight positions of the samples can now be discussed in the order of increasing distance from their ionizing stars.

5.1.1. Samples Ionized by θ^1 Ori C

Θ for Sample A and B must be negative; otherwise, samples this close to θ^1 Ori C would not be ionized by direct radiation. Sample A would have $Z = -0.09$ pc and Sample B would have $Z = -0.13$ pc, the latter being comparable to the sub- θ^1 Ori C distance and the former sample being 0.04 pc closer to the observer. This agrees with the rapid increase of surface brightness west of θ^1 Ori C.

If the values of Θ for Samples C, F, and G were negative, then this would imply that the nebula is basically flat to the southeast from θ^1 Ori C. This is incompatible with the fact that the surface brightness in $H\alpha$ is nearly constant (which requires that the surface curve toward the observer because the ionizing flux from θ^1 Ori C falls with r). These samples must have positive values of Θ and Z , which indicates that the MIF surface has risen about 0.3 pc from the sub- θ^1 Ori C point over a distance of about 0.2 pc.

Sample E certainly has positive values since it lies southeast of the Bright Bar, with this region being 0.44 pc toward the observer from the plane of θ^1 Ori C and 0.6 pc above the sub- θ^1 Ori C point at a distance of 0.3 pc. Samples C, E, F, and G taken together indicate that the MIF rises steeply toward the observer to the southeast from θ^1 Ori C, which is the direction of the Bright Bar.

Table 2
Positions of Samples with Respect to Their Ionizing Stars

Sample	Ionizing Star	r_{sky} (pc)	r/r_{sky}	Θ	Z (pc)
A	θ^1 Ori C	0.054	1.77	-60°	-0.094
B	θ^1 Ori C	0.028	4.37	-78°	-0.134
C	θ^1 Ori C	0.189	1.134	28°	0.101
D	θ^2 Ori A	0.060	2.81	-72°	-0.180
E	θ^1 Ori C	0.254	1.78	60°	0.444
F	θ^1 Ori C	0.179	1.134	28°	0.096
G	θ^1 Ori C	0.139	1.53	55°	0.196

The region outside of the Bright Bar changes at about the position of the HH 203 and HH 204 shocks. To the northeast of these shocks $F(\text{N II})/F(\text{H}\alpha)$ is high and $F(\text{O III})/F(\text{H}\beta)$ is low. This is reversed to the southwest of these shocks and agrees with our Sample D and Sample E values. This transition is illustrated well in an $F(\text{N II})/F(\text{H}\alpha)$ image, such as Figure 24 of O'Dell et al. (2015).

The southwest high-ionization region is marked by numerous crenellations lying immediately southeast of the Bright Bar. These features have sharp boundaries and are likely to be shocks driven by a series of outflows from the Orion-S star formation region (O'Dell et al. 2015).

If these shocks and earlier precursors are contributing to ionization in the region around Sample E, we can then accept that the sample does not give consistent results for the distance correction for either for θ^1 Ori C or θ^2 Ori A. It may be that θ^2 Ori A plays a secondary role in ionizing this region.

5.1.2. Samples Ionized by θ^2 Ori A

The region between Sample D and θ^2 Ori A is relatively smooth and structure-free (except for the overlying shocks of HH 203 and HH 204). This means that the sample must lie away from the observer relative to the plane of θ^2 Ori A. The progression of decreasing displacements of the neon line ratios in going from Sample I4 to Sample I1 indicates that Sample I4 is well beyond θ^2 Ori A, but by position I1 the ionization front is close to the same plane as θ^2 Ori A.

This result is quite different from that in Rubin et al. (2011), who argue in a qualitative way that all of the region along their samples (to the southeast, beyond the Bright Bar) is ionized by θ^1 Ori C, under the assumption that θ^2 Ori A is incapable of creating local $[\text{Ne III}]$ emission. Our Cloudy models indicate that this is not the case, and our examination of the variations of the $F(\text{[Ne III]})/F(\text{[Ne II]})$ ratio firmly establishes that θ^2 Ori A is the dominant ionizing star in this region.

In Section 2.4 we established that Sample E, which lies outside the Bright Bar, is dominantly ionized by θ^1 Ori C. However, this identification was good only for the $F(\text{[O III]})/F(\text{H}\beta)$ and $F(\text{He I})/F(\text{H}\alpha)$ ratios with a surprisingly large distance correction, while $F(\text{N II})/F(\text{H}\alpha)$ indicated that no distance correction was required. In the comparison with line ratios predicated on ionization by θ^2 Ori A, the observed points were close to T_{star} for that star, but always slightly hotter. It is most likely that Sample E is ionized by both θ^1 Ori C and θ^2 Ori A.

The results of this discussion are presented in Table 2.

6. Conclusions

We have confirmed the usual assumption that the bright portion of the Orion Nebula lying within the boundary of the

southeast Bright Bar is dominantly ionized by the hottest star in the Trapezium grouping (θ^1 Ori C). With the exception of a single region (Sample E), all of the regions lying southeast of the Bright Bar are dominantly ionized by the much cooler isolated star θ^2 Ori A. Sample E is most likely ionized by both stars. This conclusion is verified by the orientation of the ionization boundaries of propyds, with the exception of one propylid (196.6–422.6) that probably lies at a small spatial distance from θ^2 Ori A.

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