Magnetic and Crystal Structures of Sr$_2$IrO$_4$: A Neutron Diffraction Study

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Magnetic and crystal structures of Sr$_2$IrO$_4$: A neutron diffraction study

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We report a single-crystal neutron diffraction study of the layered Sr$_2$IrO$_4$. This work unambiguously determines the magnetic structure of the system and reveals that the spin orientation rigidly tracks the staggered rotation of the IrO$_6$ octahedra in Sr$_2$IrO$_4$. The long-range antiferromagnetic order has a canted spin configuration with an ordered moment of 0.208(3) $\mu_B$/Ir site within the basal plane; a detailed examination of the spin canting yields 0.202(3) and 0.049(2) $\mu_B$/site for the a axis and the b axis, respectively. It is intriguing that forbidden nuclear reflections of space group $I4_1/acd$ are also observed in a wide temperature range from 4 K to 600 K, which suggests a reduced crystal structure symmetry. This neutron-scattering work provides a direct, well-refined experimental characterization of the magnetic and crystal structures that are crucial to the understanding of the unconventional magnetism exhibited in this unusual magnetic insulator.

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The 5$d$-based iridates have continuously provided a fertile playground for the studies of novel physics driven by the spin-orbit interaction (SOI). It is believed that SOI (0.4–1 eV), which is proportional to $Z^3$ ($Z$ is the atomic number), plays a critical role in the iridates and rigorously competes with other relevant energies, particularly the on-site Coulomb interaction $U$ (0.4–2.5 eV), which is significantly reduced because of the extended nature of the 5$d$ orbitals. A new balance between the competing energies is, therefore, established in the iridates and drives exotic quantum phases. Recent experimental observations and theoretical proposals for the iridates have included the following: the $J_{\text{eff}} = 1/2$ Mott state, superconductivity, a correlated topological insulator with large gaps, spin liquid in a hyperkagome structure, Weyl semimetal with Fermi arcs, the Kitaev mode, and three-dimensional (3D) spin liquid with Fermionic spinons.

Among all the iridates studied, the single-layer Sr$_2$IrO$_4$ has been subjected to the most extensive investigations due to its structural and electronic similarities to the undoped high-$T_C$ cuprates such as La$_2$CuO$_4$. This magnetic insulator was proposed to be an effective $J_{\text{eff}} = 1/2$ Mott-Hubbard state arising from the SOI. Although the insulating ground state has been established by angle-resolved photoemission spectroscopy and resonant x-ray scattering (RXS) measurements, some critical insights into the crystal and magnetic structures remain conspicuously elusive. For example, the strong SOI limit $J_{\text{eff}} = 1/2$ ground-state scenario has been recently challenged by x-ray absorption spectroscopy, time-resolved optical studies, and theory. The nature of the weak ferromagnetism arising from the canted antiferromagnetic (AFM) order is not fully characterized experimentally. It is primarily due to the lack of large single crystals and the strong absorbing cross section of the Ir ions that prevent a comprehensive neutron study. Here we report the results of a neutron diffraction investigation of single-crystal Sr$_2$IrO$_4$. The central findings of this work are the following: (1) The magnetic and crystal structures are completely determined; (2) the system undergoes an antiferromagnetic transition at 224(2) K with an ordered moment of 0.208(3) $\mu_B$/Ir site and a canted spin configuration within the basal plane; and (3) the spin orientation is intimately associated with the rotation of the IrO$_6$ octahedra, which results in 0.202(3) and 0.049(2) $\mu_B$/Ir site for the a axis and the b axis, respectively. In addition, nuclear reflections incompatible with the previously reported space group (SG) are observed and indicate a possible lowering of the structural symmetry. The Sr$_2$IrO$_4$ single crystal studied (2 × 2 × 1 mm$^3$, mass = 8 mg) was grown using self-flux techniques. Because the iridium is highly neutron absorbing, the equal-dimensional shaped crystal simplifies the necessary absorption correction. The neutron diffraction measurements were carried out at the HB1A, HB1 triple axis spectrometers, and the HB3A four circle diffractometer at the High Flux Isotope Reactor at the Oak Ridge National Laboratory. For the measurements using triple axis spectrometers, the crystal was aligned in the (h,0,1), (h,h,1), (0,k,l) and other scattering planes to probe various magnetic reflections. A closed-cycle refrigerator and high temperature furnace were employed to monitor the $T$ dependence of the magnetic and nuclear reflections. Sr$_2$IrO$_4$ was reported to crystallize in a tetragonal structure (SG $I4_1/acd$, No. 142) with $a = b = 5.484$ Å and $c = 25.83$ Å at 4 K. With reflection conditions compliant with the $I4_1/acd$ symmetry, we have collected 137 nuclear reflections of Sr$_2$IrO$_4$ using HB3A for structure refinements. The most prominent features of the crystal structure are the elongation of the IrO$_6$ octahedra along the c axis (2.055 Å for the out-of-plane distance compared to 1.981 Å in-plane), and the rotation of the octahedra with respect to the c axis about 11.8(1)$^\circ$ at 4 K. This leads to a $\sqrt{2} \times \sqrt{2}$ expansion of unit cell in the basal plane compared to the higher symmetry Sr$_2$RuO$_4$ [Fig. 1(a)].

The antitranslation in combination with the body centering dictates a (1,1,1) magnetic propagation wave-vector, as discussed previously. Figure 2(a) displays the $T$ dependence of the Bragg intensity ($I_B \propto |M|^2$, $M$ is the order parameter) of the magnetic reflection (1,0,2). The intensity vanishes around $T_N = 224(2)$ K and is consistent with the...
magnetization measurement.\textsuperscript{17} Fitting the order parameter to the power-law scaling function $I_B \sim |\psi|^\beta$, where $\tau = 1 - T/T_N$ is the reduced temperature, leads to the critical exponent $\beta = 0.18(1)$. It apparently deviates from the $\beta = 0.325$ expected for a 3D Heisenberg spin system. Figures 2(b) and 2(c) illustrate the wave-vector scans within and perpendicular to the basal plane at several temperatures. In both cases, the lineshape of the magnetic scattering evolves into a Gaussian profile below $T_N$, signaling the formation of the long-range magnetic order. Our observation is in accord with the RXS studies indicating that a short-range Heisenberg spin fluctuation occurs only in a paramagnetic state.\textsuperscript{22} A quantitative characterization of the magnetic structure and moment size of the Ir\textsuperscript{4+} ions can be obtained by a comprehensive survey of the magnetic reflections in conjunction with the model calculation. Figure 3 shows the neutron diffraction scans at selected reflections. The disappearance of the scattering above $T_N$ and decrease in intensity at large momentum transfer indicate their magnetic nature. Differing from the early RXS studies where the magnetic reflections are present only at $(0,1,4n + 2)$ and $(1,0,4n)$,\textsuperscript{3} our neutron diffraction shows additional Bragg peaks at the $(0,1,4n)$ and $(1,0,4n + 2)$ positions. The nearly identical intensity at equivalent wave-vectors $(1,0,2)$ and $(0,1,2)$ indicates the crystal has equally populated magnetic domains. Note that the structural refinement with the same sample cannot determine whether the system is structurally twinned.\textsuperscript{23} The presence of both types of reflections strongly suggests that they originate from the twinned crystallographic domains. According to the Landau theory, the symmetry properties of the magnetic structure can be described by only one irreducible representation (IR). With Ir ions located at the 8a Wyckoff positions for the SG $I4_1/acd$ and the propagation wave-vector $q_{M} = (1,1,1)$, the magnetic representation can be decomposed into $\Gamma_{mag} = 2\Gamma_1^2 + 2\Gamma_2^2 + 2\Gamma_3^2 + 2\Gamma_4^2$, where $\Gamma_1$, $\Gamma_2$ are the two-dimensional IRs with basis vectors lying in the $ab$ plane and $\Gamma_3$, $\Gamma_4$ are the one-dimensional IRs with moments pointing parallel to the $c$ axis. Since the magnetic susceptibility suggests that the spin easy axis lies in the basal plane,\textsuperscript{17} we exclude spin configurations associated with $\Gamma_3$ and $\Gamma_4$ in the analysis. Table I lists the basis vectors of IRs $\Gamma_1$ and $\Gamma_2$. In particular, the spin structure based on linear combination of $\psi(2)$ and $\psi(3)$ of $\Gamma_1$ has a $(+-+--)$ configuration along the $a$ axis (or the M4 structure described in Refs. 16 and 21) and $(+++-)$ along the $b$ axis for the labeled Ir ions in Fig. 1. In contrast, the linear combination of $\psi(5)$ and $\psi(8)$ in $\Gamma_3$ gives $(+-+-)$ along the $a$ axis and $(+++-)$ along the $b$ axis (the M3 configuration). These spin structures derived from representation analysis using BasIreps program\textsuperscript{24} are in accord with the results from previous neutron powder diffraction.\textsuperscript{21} Table II compares the expected intensities for the two relevant spin models and the experimental observations. The M4 and M3 configurations each have distinct distributions of magnetic scattering intensities.\textsuperscript{26} For example, the collinear structure with $a$ axis $(+-+--)$ components produces the strongest scattering at the $(0,1,2)$ reflection and gives zero intensity at the $(1,0,0)$ Bragg point. However, the $(+++-)$ collinear state associated with the M3 configuration will generate the strongest scattering at the $(1,0,0)$ peak, which is not observed experimentally. The neutron diffraction results shown in Table II clearly support the M4 spin configuration and confirm the previous neutron diffraction work on the powder sample.\textsuperscript{21} To test whether there are additional canted moments along the $b$ axis with the $(+++-)$ configuration within $\Gamma_1$, we probed the scattering at the expected $(0,0,2n + 1)$ reflections. Figures 3(d) and 3(e) display the scans at the $(0,0,3)$
TABLE II. Comparison of observed and calculated magnetic intensities from two symmetry compatible spin models. To get the scale factor, separate sets of nuclear reflections were collected at HB3A with incident neutron wavelength of 1.5424 and 1.003 Å, respectively. Additional 37 nuclear reflections were collected at HB1A for intensity normalization.25

<table>
<thead>
<tr>
<th>Reflection</th>
<th>Observation</th>
<th>M4 model</th>
<th>M3 model</th>
</tr>
</thead>
<tbody>
<tr>
<td>(0,0,3)</td>
<td>0.26 ± 0.03</td>
<td>0.25</td>
<td>0.25</td>
</tr>
<tr>
<td>(0,0,5)</td>
<td>0.20 ± 0.03</td>
<td>0.20</td>
<td>0.20</td>
</tr>
<tr>
<td>(1,1,1)</td>
<td>0.08 ± 0.07</td>
<td>0.08</td>
<td>0.08</td>
</tr>
<tr>
<td>(0,1,2)</td>
<td>6.80 ± 0.17</td>
<td>6.99</td>
<td>1.05</td>
</tr>
<tr>
<td>(0,1,6)</td>
<td>4.72 ± 0.32</td>
<td>4.50</td>
<td>2.73</td>
</tr>
<tr>
<td>(1,0,2)</td>
<td>6.99 ± 0.18</td>
<td>6.99</td>
<td>1.05</td>
</tr>
<tr>
<td>(1,0,4)</td>
<td>2.33 ± 0.22</td>
<td>2.48</td>
<td>5.81</td>
</tr>
<tr>
<td>(1,0,6)</td>
<td>4.72 ± 0.32</td>
<td>4.51</td>
<td>2.73</td>
</tr>
<tr>
<td>(1,0,8)</td>
<td>2.56 ± 0.32</td>
<td>2.28</td>
<td>3.01</td>
</tr>
<tr>
<td>(1,0,14)</td>
<td>0.88 ± 0.21</td>
<td>0.54</td>
<td>0.48</td>
</tr>
<tr>
<td>(1,0,16)</td>
<td>0.18 ± 0.09</td>
<td>0.24</td>
<td>0.26</td>
</tr>
<tr>
<td>(1,2,0)</td>
<td>1.53 ± 0.36</td>
<td>1.76</td>
<td>0.43</td>
</tr>
<tr>
<td>(1,2,4)</td>
<td>1.14 ± 0.23</td>
<td>1.46</td>
<td>0.52</td>
</tr>
<tr>
<td>(1,2,8)</td>
<td>0.85 ± 0.12</td>
<td>0.82</td>
<td>0.45</td>
</tr>
</tbody>
</table>

and (0,0,5) Bragg peaks. Although much weaker, the magnetic scattering is clearly present at low $T$ and confirms the staggered AFM order propagating along the c axis. A total of 14 magnetic reflections combined with 137 nuclear reflections allow an accurate determination of the spin structure and the associated moment. Using the M4 spin model and the magnetic form factor for Ir$^{4+}$,27 we have obtained $m_a = 0.202(3)(μ_B$ along the a axis and $m_b = 0.048(2)(μ_B$ along the b axis, yielding a total moment of $0.208(3)(μ_B/Ir^{4+}$ site. This value is smaller than 0.36(6) $μ_B$ from a recent single crystal neutron-scattering study28 but quite consistent with the powder neutron diffraction results in which the upper limit of the moment does not exceed 0.29(4) $μ_B$.21 The magnetic configuration in Figs. 1(b)–1(d) show that spins projected along the b axis have a staggered $↓↑↑↓$ pattern along the c axis, with Ir spins deviating 13(1)$°$ away from the a axis [see Fig. 1(d)]. This spin canting rigidly tracks the staggered octahedral rotation, as illustrated in a previous RXS study.1 This remarkable correlation proves the existence of strong magnetoelastic coupling in the iridate, which is also suggested in experimental studies of transport and magnetic properties of the system.4,29

Theoretically, the spin Hamiltonian in the strong SOI limit includes the isotropic coupling ($J$) and the Dzyaloshinskii-Moriya interaction ($D$) caused by the lattice distortion.11

FIG. 2. (Color online) (a) The $T$-dependence of the magnetic (1,0,2) reflection. Inset shows the intensity versus the reduced temperature ($t = |1 − T/T_N|$) in logarithmic scale. The wave-vector scan along (b) the [H,0,0] and (c) the [0,0,L] directions for the (1,0,2) peak at selected temperatures that probe the in-plane and out-of-plane correlations. (d) Similar wave-vector scans for the magnetic (1,1,1) reflection above and below $T_N$. Note that the counting time is 10 times compared to those of the strong (1,0,2) peak.
The spin canting is governed by the ratio of $D$ and $J$ and is solely determined by the lattice distortion. This explains the relatively large spin canting in the $5d$ system compared to that in La$_2$CuO$_4$ where SOI is insignificant (SOI the relatively large spin canting in the 5 $d$ system is solely determined by the lattice distortion. This explains the (0,0,5) magnetic reflections. The weaker (0,0,2n + 1) are measured with much longer counting time. 

The intensity continuously decreases on warming and shows no sign of transition to 600 K. The reduction in intensity cannot be accounted by the thermal vibration of the elements (Debye-Waller factor). The lack of anomaly near $T_N$ is also consistent with the transport, thermodynamic, and optical conductivity studies. Such observation of reduced structural symmetry that persists at a much higher temperature than $T_N$, implies the formation of a crystallographic template for the low-$T$ spin structure that changes the tetragonal symmetry. This observation is certainly intriguing and the origin of it remains to be understood.

It is established that the magnetic and electronic properties are highly susceptible to slight impurity doping for Sr, Ir, or oxygen. For example, doping Mn results in a spin-flop transition with moments aligning along the $c$ axis. The remaining $J_{\text{eff}} = 1/2$ state revealed by RXS measurement suggests its robustness against the alternation of spin structure.

FIG. 3. (Color online) Selected rocking scans at 4 K and 250 K for the (a) (1,0,2), (b) (1,0,4), (c) (1,0,6), (d) (1,0,8), (e) (0,0,3), and (f) (0,0,5) magnetic reflections. The weaker (0,0,2n + 1) are measured with much longer counting time.
FIG. 4. (Color online) (a) The $T$-dependence of the structural (0,1,1) reflection. Open circles are the peak intensity and solid squares the integrated intensity. Dashed line is the background derived from the Gaussian fit to the rocking scan. (b) The rocking scans of the (0,1,1) peak at $T = 60$, 150, 280, 390, and 500 K. The wave-vector scans along the [1,0,0] direction for (c) the (1,0,1) and (d) the (1,0,5) reflections at 5 K and 250 K.

On the other hand, replacing Ir with isovalent Rh$^{4+}$ leads to a rich phase diagram of metal-insulator transition tuned by SOI. The transition was explained by the effective reduction of the splitting between the $J_{\text{eff}} = 1/2$ and $J_{\text{eff}} = 3/2$ bands due to the reduced SOI; this in turn alters the relative strength of the SOI and the crystal electric field (CEF) that dictates the magnetic state. This notion is also consistent with a recent theoretical proposal that the change of CEF associated with the underlying structure could be critical to determine the magnetic ground states. The present single-crystal neutron diffraction unambiguously determines the magnetic structure and proves the rigid coupling of the spin canting with the rotation of the IrO$_6$ octahedra. These findings finally fill the longstanding gap in our understanding of the magnetic properties in Sr$_2$IrO$_4$, an archetype of the $J_{\text{eff}} = 1/2$ insulators.

We thank Q. Huang, S. Lovesey, D. Khalyavin, and G. Khaliullin for invaluable discussions. Research at ORNL's High Flux Isotope Reactor was sponsored by the Scientific User Facilities Division, Office of Basic Energy Sciences, US Department of Energy. The work at University of Kentucky was supported by NSF through Grants No. DMR-0856234 and No. EPS-0814194.

18 The absorption correction factor has been numerically applied to the collected reflections based on the sample shape. The mean transmission ($T$) is 0.92 with the minimum at the (0, 0, 8) reflection ($T = 0.79$) and the maximum at the (2, 1, 1) peak ($T = 0.95$).
23 Our neutron diffraction measurement typically has beam size of 5 × 5 mm at the sample position and is larger than the 10 ∼ 100 μm used in the RXS experiment that is comparable with magnetic domain size. See, for example, S. Bosegga, R. Springell, H. C. Walker, A. T. Boothroyd, D. Prabhakaran, D. Wermille, L. Bouchenoire, S. P. Collins, and D. F. McMorrow, Phys. Rev. B 85, 184432 (2012).
25 The nuclear reflections are collected using two-axis mode, while the magnetic reflections are collected using three-axis mode to improve the signal-to-noise ratio. The corresponding integrated intensities are corrected using the method described by R. Pynn, Acta Cryst. B 31, 2555 (1975).
26 Because of the tetragonal symmetry of the crystal structure, the linear combination of $\psi(1)$ and $\psi(4)$ within the same $\Gamma$ IR can also describe the observed intensities. For simplicity, we use $\psi(2)$ and $\psi(3)$ to be consistent with the spin configuration reported in Ref. 21.